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VARIATIONAL STABILITY AND LOCAL RIGIDITY OF EINSTEIN METRICS

By

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Abstract. Certain curvature conditions for variational stability of Einstein metrics are given. The argument of Besson, Courtois and Gallot, developed in [2], is improved in terms of the Weyl curvature operator and the scalar curvature. For compact Kähler-Einstein manifolds (M, g, J) we can show that the infinitesimal rigidity of complex structures J is equivalent to the variational stability of Einstein metrics g.

1. Introduction

An Einstein metric g on a compact manifold M is called *locally rigid*, if it gives an isolated point [g] in the moduli space of Einstein metrics on M. Further, an Einstein metric g is called *variationally stable*, when the quadratic form associated to the second variation of the total scalar curvature functional $S(g') = \operatorname{Vol}(g')^{2/n-1} \int_M s_{g'} dv_{g'}$ is positive definite; i.e., there exists $\lambda > 0$ such that

$$-S''_{q}(h,h) \ge \lambda \|h\|^2$$

for any $h \in \Gamma(M; S^2(M))$ satisfying $\operatorname{tr}_g h = 0$ and $\delta_g h = 0$. See Definition 4.63 in [1].

For the precise definition of local rigidity refer to [10], where (and also in [1]) the terminology *rigidity* is used. Notice also that by the result of [9] any variationally stable Einstein metric is locally rigid, since any variationally stable Einstein metric admits no non-trivial infinitesimal Einstein deformation.

Several examples of locally rigid Einstein metrics are given in [1]; the *n*-sphere S^n with the standard metric, the complex projective space $\mathbb{C}P^n$ with the Fubini-Study metric and a compact Einstein manifold of negative sectional curvature.

For Einstein manifolds of negative sectional curvature Besson, Courtois and

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Gallot considered in [2] the functional $g' \mapsto K(g') = \int_X |s_{g'}|^{n/2} dv_{g'}$ which is closely related to the functional S = S(g') and obtained the following variational stability theorem.

THEOREM (Besson, Courtois and Gallot [2]). Let (X, g) be a compact, connected oriented n-manifold X with an Einstein metric g of negative scalar curvature. If g is a metric of negative sectional curvature or a locally symmetric metric of non-compact type, then $K : \Sigma \to \mathbf{R}$ is locally minimal at g' = g. Here Σ denotes the set of all smooth metrics g' with vol(g') = 1 and of constant scalar curvature. Namely, such an Einstein metric is variationally stable and hence locally rigid.

The local rigidity of Einstein manifolds of negative curveture is also stated in [1], Corollary 12.73.

The following theorem, known as a global rigidity theorem, states that the moduli of Einstein metrics on hyperbolic manifolds consists of a single point.

THEOREM (Besson, Courtois and Gallot [3], LeBrun [13]). Let X be a compact, connected oriented 4-manifold admitting a real or complex hyperbolic metric g_0 . Then any Einstein metric g on X is homothetic to g_0 up to diffeomorphisms of X.

The aim of this article is to relax the strictly negative curvature condition of the local rigidity of Einstein metrics in Theorem (Besson, Courtois and Gallot [2]). Since the Ricci tensor is a multiple of the metric, the Riemannian curvature tensor R is expressed as the sum of the Weyl conformal curvature part W and the scalar curvature part and we obtain the following variational stability theorems.

THEOREM 1. Let (M, g) be a compact, connected oriented Einstein n-manifold with $s_g < 0$. If

$$\sup_{x\in M} w(x) + \frac{1}{n(n-1)}s_g < 0,$$

then g is variationally stable and then locally rigid. Here w(x) denotes the largest eigenvalue of the Weyl curvature operator $W : \Lambda^2(M) \longrightarrow \Lambda^2(M)$ at x.

THEOREM 2. Let (M, g) be a compact, connected oriented Einstein 4-manifold. If $s_g < 0$ and

$$\sup_{x \in M} \{w^+(x) + w^-(x)\} + \frac{s_g}{6} < 0,$$

then q is variationally stable and hence locally rigid.

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Here $w^+ = w^+(x)$ and $w^- = w^-(x)$ denote the largest eigenvalue of the (anti-)self-dual Weyl curvature operators $W^{\pm} : \Lambda^2_{\pm}(M) \longrightarrow \Lambda^2_{\pm}(M)$ at x, respectively. Here $\Lambda^2_{\pm}(M)$ are the bundles of self-dual, or anti-self-dual 2-forms.

COROLLARY 1. Let (M,g) be a compact, connected oriented Einstein 4-manifold. Suppose that $s_g < 0$ and g is self-dual, i.e., $W^- = 0$ identically. If $\sup_M w^+ + \frac{s_g}{6} < 0$, then g is variationally stable and then locally rigid.

This corollary supports strongly the following conjecture; a compact, oriented self-dual Einstein 4-manifold (M, g) with $s_g < 0$ must be real hyperbolic or complex hyperbolic.

Now, let (M,g) be a compact Kähler-Einstein manifold of real dimension four with $s_g < 0$. Then from Proposition 2, [4] the largest eigenvalue of W^+ is exactly $w^+ = -\frac{1}{12}s_g$.

COROLLARY 2. Let (M, g) be a compact, connected Kähler-Einstein manifold of real dimension four. Suppose that $s_g < 0$ and $\sup_M w^- + \frac{s_g}{12} < 0$. Then g is variationally stable and hence locally rigid.

For compact Kähler-Einstein manifolds of $s_g < 0$ and of dimension n for arbitrary n we are able to present a similar stability theorem in terms of the Bochner curvature tensor(indeed, as Theorem 4.1 in §4). However, due to Calabi-Aubin-Yau's argument, deformation of Einstein metrics on a complex manifold is directly related to the deformation of complex structures. By the result of Koiso, if a compact Kähler-Einstein manifold of $s_g < 0$ is variationally stable, then the first Kodaira-Spencer cohomology group vanishes, i.e., $H^1(M; \Theta) = 0$, that is, the complex structure is infinitesimally rigid(see [8]). See for this statement [11] and Proposition 12.98, [1]. We derive the following theorem which is the converse implication.

THEOREM 3. Let (M, g, J) be a compact Kähler-Einstein manifold of $s_g < 0$. If $H^1(M; \Theta) = 0$, then g is variationally stable.

2. The functionals K and S

Let \mathcal{R} be the space of all smooth metrics on a compact oriented *n*-manifold M. Define functionals $K, S : \mathcal{R} \longrightarrow \mathbf{R}$ by

$$K(g') = \int_M |s_{g'}|^{n/2} dv_{g'},$$

$$S(g') = \operatorname{vol}(g')^{2/n-1} \int_M s_{g'} dv_{g'}, \ g' \in \mathcal{R}.$$

These functionals K, S are significantly important for studying geometry of Riemannian manifolds and appear for instance in the Yamabe problem associated to scalar curvature (see [12], and [15]).

THEOREM 2.1 (Besson, Courtois and Gallot [2]). Let g be a smooth metric on a compact oriented n-manifold M having negative constant scalar curvature. If g' is a smooth metric and is conformal to g, then

 $K(g') \ge K(g),$

(1)

and the equality holds if and only if g' = cg for some positive constant c.

Remark that the same statement is also given in [15]. In four dimension the Seiberg-Witten theory enables to estimate for certain 4-manifolds the absolute minimum of K, the square L^2 -functional of scalar curvature, in terms of the topological invariants, as

$$K(g') \ge 32\pi^2(2\chi(M) + 3\tau(M)),$$

for any smooth metric g' and the equality holds when the 4-manifold is Kähler-Einstein. See for this [5], [14].

Denote by \mathcal{R}^{vol} the space of smooth metrics of unit volume,

$$\mathcal{R}^{vol} = \{ g' \in \mathcal{R} \mid \operatorname{vol}(g') = 1 \},\$$

and by \mathcal{R}^{scal} the space of smooth metrics of constant scalar curvature,

$$\mathcal{R}^{scal} = \{ g' \in \mathcal{R} \mid s_a \text{ is constant } \},\$$

and restrict the functional K to $\Sigma = \mathcal{R}^{vol} \cap \mathcal{R}^{scal}$, the space of metrics of unit volume having constant scalar curvature. Then

$$K(g') = |s_{g'}|^{n/2} = (-s_{g'})^{n/2} = \{-S(g')\}^{n/2}.$$

Therefore, $K|_{\Sigma}$ is locally minimal at g' = g if and only if $S|_{\Sigma}$ is locally maximal at g' = g.

The functional S is invariant under the action of the group $\text{Diff}^+(M)$ of orientation preserving diffeomorphisms of M. The space \mathcal{R}^{vol} is also invariant under its action. Notice that for any $g \in \mathcal{R}^{vol}$ there exists a slice $\mathcal{R}^{vol}_{a} \subset$

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 \mathcal{R}^{vol} transversal to the action of $\text{Diff}^+(M)$ such that $g \in \mathcal{R}^{vol}_o$ and the subset $\{\phi^*(g') \mid g' \in \mathcal{R}^{vol}_o, \phi \in \text{Diff}^+(M)\}$ contains a neighborhood of g in \mathcal{R}^{vol} .

Then the tangent space at g to \mathcal{R}^{vol} is given by the direct sum

$$T_g \mathcal{R}^{vol} = \mathcal{L}_{\mathfrak{X}(M)} \oplus T_g \mathcal{R}_o^{vol}$$

where

$$\mathcal{L}_{\mathfrak{X}(M)} = \{ h \in \Gamma(M; S^2M) \mid h = \mathcal{L}_X g \text{ for } X \in \mathfrak{X}(M) \}$$

is the infinite dimensional vector space tangent to the $\text{Diff}^+(M)$ -orbit through g and

$$T_g \mathcal{R}_o^{vol} = \{ h \in \Gamma(M; S^2 M) \mid \int_M \operatorname{tr}_g h \, dv_g = 0, \, \delta_g h = 0 \, \}$$

is the tangent space at g to \mathcal{R}_{o}^{vol} .

Any 2-symmetric tensor h is written as $h = h_0 + fg$, where h_0 is trace free and $f \in C^{\infty}(M)$. The space $T_g \mathcal{R}_o^{vol}$ then decomposes into the subspaces $\Phi \oplus C_o^{\infty}(M)g$, where

$$\Phi = \{h \mid \operatorname{tr}_{q} h = 0, \, \delta_{q} h = 0\}$$

and

$$C_o^{\infty}(M)g = \{h = fg \, | \, f \in C^{\infty}(M), \, \int_M f dv_g = 0 \, \},$$

the tangent space to the conformal deformations $\{e^f g \in \mathcal{R}^{vol} \mid f \in C^{\infty}(M)\}$ of the metric g.

The first variation formula of $S : \mathcal{R} \longrightarrow \mathbf{R}$ is

$$-S'_g(h) = \int_M \left(r_g - rac{s_g}{2}g, h
ight)_g dv_g,$$

where r_g is the Ricci tensor and $g(t) : (-\epsilon, \epsilon) \longrightarrow \mathcal{R}$ is a curve with g(0) = gand $\frac{d}{dt} g(t)|_{t=0} = h$. This formula is derived from the following formulae([10], [1]);

$$\begin{split} & \left(\frac{d}{dt}s_{g(t)}\right)|_{t=0} = \Delta_g(\mathrm{tr}_g h) + \delta_g(\delta_g h) - (r_g, h)_g \\ & \left(\frac{d}{dt}dv_{g(t)}\right)_{|t=0} = \frac{1}{2}\mathrm{tr}_g h \, dv_g \end{split}$$

Therefore, we have

PROPOSITION 2.1. Let g be a smooth metric of negative constant scalar curvature s_g . Then, the following are equivalent.

- (i) g is Einstein,
- (ii) g is critical for the total scalar curvature functional $S: \Sigma \longrightarrow \mathbf{R}$.

This proposition, a well known result of D. Hilbert, appears in [1] as Proposition 4.47. The proof of this proposition is easy so that we omit it.

3. The quadratic form Q

Let g be an Einstein metric on a compact manifold M with $s_q < 0$.

PROPOSITION 3.1 (Koiso [9], Theorems 2.4, 2.5 and Besse [1], Theorem 4.60). The second variation formula of $S: \Sigma \cap \mathcal{R}_o^{vol} \longrightarrow \mathbf{R}$ at g is given by

$$-2S''_{g}(h,h) = \int_{M} \{|Dh|^{2} - 2(\hat{R}(h),h)\} dv_{g},$$
⁽²⁾

where $\hat{R}: S_0^2(M) \longrightarrow S_0^2(M)$, $S_0^2(M) = \{h \in S^2(M) | \operatorname{tr}_g h = 0\}$, is an endomorphism defined by

$$\hat{R}(h)_{ij} = \sum_{k,\ell} R_{ikj\ell} h_{k\ell}$$

in terms of an orthonormal basis. Here $R = (R_{ikj\ell})$ is the Riemannian curvature tensor of (M, g).

We remark here that the tangent space at g to $\Sigma \cap \mathcal{R}_o^{vol}$ is given, as noticed in p.131, [1], by

$$T_g(\Sigma \cap \mathcal{R}_o^{vol}) = \{h \in \Gamma(M; S^2M) \mid \operatorname{tr}_g h = 0, \, \delta_g h = 0\}.$$

In fact, take a one-parameter family g(t) in $\Sigma \cap \mathcal{R}_o^{vol}$ with g(0) = g. Then $h = \frac{d}{dt}g(t)_{|t=0}$ satisfies $\delta_g h = 0$ and $\int_M \operatorname{tr}_g h dv_g = 0$. Also from the scalar curvature constant condition we have $\Delta_{g(t)}s_{g(t)} = 0$ and by differentiating this by t

$$\Delta_g(\Delta_g(\mathrm{tr}_g h) + \delta_g(\delta_g h) - \frac{s_g}{n}\mathrm{tr}_g h) = 0.$$

Since $\delta_g h = 0$, $\Delta_g(\operatorname{tr}_g h) - \frac{s_g}{n} \operatorname{tr}_g h$ must be constant. However, this constant must be further zero, because $\int_M \operatorname{tr}_g h dv_g = 0$. Then $\operatorname{tr}_g h$ is an eigenfunction of negative eigenvalue $\frac{s_g}{n}$. This implies that h is trace free.

Now we will investigate the second variation formula (2). We define the quadratic form Q associated with this formula.

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The rough Laplacian D^*D for $h \in \Gamma(M; S^2_0(M))$ can be written as

$$D^*Dh = (\delta^D_g d^D_g + d^D_g \delta^D_g)h - h \circ r + \hat{R}(h), \qquad (3)$$

where

$$(h \circ r)_{ij} = \sum_{k} h_{ik} r_{kj}.$$

We will follow the argument given by [2].

Take an arbitrary $\alpha \in (0,1)$. Then, since $\delta_g h = 0$, it holds from the above formula that

$$-2S_{g}''(h,h) = \alpha \int_{M} |Dh|^{2} dv_{g} + (1-\alpha) \int_{M} |d_{g}^{D}h|^{2} dv_{g}$$
$$-2\alpha \int_{M} (\hat{R}(h),h) dv_{g} + (1-\alpha) \int_{M} Q(h,h) dv_{g}.$$
(4)

Here the quadratic form $Q = Q_x$ at $x \in M$ is defined by

$$Q_x(h,h) = -\{(h \circ r, h)_g(x) + (\hat{R}(h), h)_g(x)\} \text{ for } h \in S_0^2(M).$$
(5)

The positive definiteness of Q gives a sufficient criterion to the variational stability of Einstein metrics in the following lemma.

LEMMA 3.1 (Besson, Courtois and Gallot [2]). If there exists a positive constant λ such that for any $x \in M$ and for any h of $\operatorname{tr}_{g} h = 0$

 $Q_x(h,h) \ge \lambda |h|^2,$

then there exists also $\lambda' > 0$ such that

$$-2S_g''(h,h) \ge \lambda' \int_M |h|_g^2 dv_g.$$
(6)

for all h.

If, in fact, we take $\alpha > 0$ sufficiently small, then we obtain (6) from (4).

We will verify our theorems and corollaries stated in Introduction. Let g be an Einstein metric of $s_g < 0$. Then, since $r = \frac{s_g}{n}g$, the Riemannian curvature tensor R_g is written as

$$R_g = W - \frac{s_g}{n(n-1)}g \bigotimes g,$$

where W is the Weyl curvature tensor, so that

$$\hat{R}(h) = \hat{W}(h) - \frac{s_g}{n(n-1)} \ h, \ h \in S_0^2(M).$$

The quadratic form $Q_x(h,h)$ is therefore

$$Q_x(h,h) = -\left\{ \frac{n-2}{n(n-1)} s_g |h|^2 + (\hat{W}(h),h) \right\}.$$

Diagonalize h as

$$h(e_i, e_i) = h_{ii}, \quad h(e_i, e_j) = h_{ij} = 0, \quad i \neq j.$$

We have then $|h|^2 = \sum_{i=1}^n h_{ii}^2$ and

$$\hat{W}(h)_{ij} = \sum_{k,\ell} W_{ikj\ell} h_{k\ell} = \sum_{k=1}^{n} W_{ikjk} h_{kk}$$

so that

$$(\hat{W}(h),h) = \sum_{i,j} (\hat{W}(h))_{ij} h_{ij} = \sum_{i,k} (W_{ikik}h_{kk})h_{ii}$$
$$= \sum_{i \neq k} W_{ikik}h_{ii}h_{kk}$$

which turns out to be

$$\frac{1}{2} \sum_{i \neq k} W_{ikik} (h_{ii} + h_{kk})^2 = \sum_{i < k} W_{ikik} (h_{ii} + h_{kk})^2,$$

since $\sum_{i} W_{ikij} = 0$. We have

$$W_{ikik} = (W(\theta^i \wedge \theta^k), \theta^i \wedge \theta^k) \le w |\theta^i \wedge \theta^k|^2 = w$$

for each $i, k, i \neq k$, where w = w(x) is the maximum eigenvalue of the self-adjoint operator $W : \Lambda^2(M) \longrightarrow \Lambda^2(M)$ at $x \in M$. Then

$$\sum_{i < k} W_{ikik} (h_{ii} + h_{kk})^2 \le w(x) \sum_{i < k} (h_{ii} + h_{kk})^2 = w \left\{ (n-1) \sum_i h_{ii}^2 + 2 \sum_{i < k} h_{ii} h_{kk} \right\}$$

Further we have from $\operatorname{tr}_g h = 0$

$$2 \sum_{i < k} h_{ii} h_{kk} = -\sum_{i} h_{ii}^2 = -|h|^2$$

from which it follows

$$(\hat{W}(h),h) \leq (n-2) w(x) |h|^2, h \in S_0^2(M)$$

and then

$$\frac{n-2}{n(n-1)}s_g|h|^2 + (\hat{W}(h),h) \leq (n-2)(w(x) + \frac{1}{n(n-1)}s_g)|h|^2$$

at each x.

Since, by the assumption,

$$\sup_{x\in M}w(x)+\frac{1}{n(n-1)}s_g<0,$$

there exists a positive constant λ such that

$$Q_x(h,h) \ge \lambda |h|^2, \ h \in S^2_0(M)$$

therefore for sufficiently small α we can use Lemma 3.1 to have

$$-2S''(h,h) \ge \lambda' \|h\|^2$$

for any $h \in \Gamma(M; S_0^2(M))$ satisfying $\delta_g h = 0$. This verifies Theorem 1.

Proof of Theorem 2. Since dim M = 4, we have $h_{ii} + h_{jj} = -(h_{kk} + h_{\ell\ell})$ for distinct i, j, k, ℓ so that

$$\sum_{i < k} W_{ikik} (h_{ii} + h_{kk})^2 = (W_{1212} + W_{3434})(h_{11} + h_{22})^2 + (W_{1313} + W_{2424})(h_{11} + h_{33})^2 + (W_{1414} + W_{2323})(h_{11} + h_{44})^2.$$

We have

$$W_{1212} + W_{3434} = (W(\theta^1 \wedge \theta^2), \theta^1 \wedge \theta^2) + (W(\theta^3 \wedge \theta^4), \theta^3 \wedge \theta^4)$$

and the similar formulae for $W_{1313} + W_{2424}$ and $W_{1414} + W_{2323}$.

Since the space of 2-forms $\Lambda^2(M)$ splits into $\Lambda^2_+ \oplus \Lambda^2_-$ equipped with the orthogonal basis

$$\begin{split} f_1^{\pm} &= \frac{1}{2}(\theta^1 \wedge \theta^2 \pm \theta^3 \wedge \theta^4), \\ f_2^{\pm} &= \frac{1}{2}(\theta^1 \wedge \theta^3 \pm \theta^4 \wedge \theta^2), \\ f_3^{\pm} &= \frac{1}{2}(\theta^1 \wedge \theta^4 \pm \theta^2 \wedge \theta^3), \end{split}$$

respectively, we have

$$\theta^{1} \wedge \theta^{2} = f_{1}^{+} + f_{1}^{-}, \quad \theta^{3} \wedge \theta^{4} = f_{1}^{+} - f_{1}^{-}$$

so that

$$W_{1212} = (W(f_1^+ + f_1^-), f_1^+ + f_1^-) = (W^+(f_1^+), f_1^+) + (W^-(f_1^-), f_1^-)$$

which coincides with W_{3434} so that

$$W_{1212} + W_{3434} = 2\{(W^+(f_1^+), f_1^+) + (W^-(f_1^-), f_1^-)\}.$$

Here we used the four dimensional Weyl curvature property, namely, $W : \Lambda^2(M) \rightarrow \Lambda^2(M)$ decomposes into $W^{\pm} : \Lambda^2_{\pm} \longrightarrow \Lambda^2_{\pm}$.

Let w^+ and w^- be the maximum eigenvalue of W^{\pm} , respectively. Then, we have

$$W_{1212} + W_{3434} \le 2 (w^+ |f_1^+|^2 + w^- |f_1^-|^2) = w^+ + w^-$$

(here $|f_i^{\pm}|^2 = \frac{1}{2}$, i = 1, 2, 3) and hence

$$\sum_{i < k} W_{ikik} (h_{ii} + h_{kk})^2 \le (w^+ + w^-) \{ (h_{11} + h_{22})^2 + (h_{11} + h_{33})^2 + (h_{11} + h_{44})^2 \}$$
$$= (w^+ + w^-) |h|^2.$$

From the assumption of Theorem 2, $\sup_{x \in M} \{w^+(x) + w^-(x)\} + \frac{1}{6}s_g < 0$. Then, it follows that for some positive constant λ

$$Q_x(h,h) \ge \lambda |h|^2, \quad h \in S^2_0(M)$$

at each point x of M. This implies from Lemma 3.1 the variational stability of our Einstein metric g.

4. Bochner curvature and Kähler-Einstein metrics

The Bochner curvature tensor $B = B_{ikj\ell}$ is defined on a Kähler manifold (M, J, g), quite similarly to the Weyl conformal curvature tensor on a real Riemannian manifold. The following stability theorem is a generalization of Corollary 2 in §1, because in four dimension B coincides with W^- .

THEOREM 4.1. Let (M, g, J) be a compact Kähler-Einstein manifold of real dimension n. Suppose $s_g < 0$. If the Bochner curvature tensor $B = B_{ikj\ell}$ satisfies

$$\sup_x \beta^+(x) + \frac{1}{n+2}s_g < 0,$$

then g is variationally stable. Here β^+ denotes the largest eigenvalue of the selfadjoint endomorphism B over the space $\{h \in S_0^2(M) \mid h \text{ is anti-J-invariant}\}.$ *Proof.* As explained in the proof of Lemma 12.94, [1], the operator $D^*D - 2\hat{R}$ preserves *J*-invariant and anti-*J*-invariant symmetric tensors. Therefore, for $h = h_1 + h_2$, where h_1 is *J*-invariant and h_2 is anti-*J*-invariant, we have

$$((D^*D - 2\hat{R})h, h)_g = ((D^*D - 2\hat{R})h_1, h_1)_g + ((D^*D - 2\hat{R})h_2, h_2)_g.$$

For J-invariant h_1 , set a 2-form ψ as $\psi(X, Y) = h_1(X, JY)$. Then by arranging the Weitzenböck formula 12.92' in [1] as

$$((D^*D - 2\hat{R})h_1)(X, Y) = -2\frac{s_g}{n}h_1(X, Y) - (\Delta\psi)(X, JY),$$

we obtain

$$((D^*D - 2\hat{R})h_1, h_1)_g = -2\frac{s_g}{n}|h_1|^2 + 2(\Delta\psi, \psi)_g$$

and by integrating

$$-2S''(h_1,h_1) = \int_M ((D^*D - 2\hat{R})h_1,h_1)_g \, dv_g \ge -2\frac{s_g}{n} \int_M |h_1|_g^2 \, dv_g,$$

since $\int_{\mathcal{M}} (\Delta \psi, \psi) dv_g \ge 0$.

On the other hand, for an anti-J-invariant h_2 we make use of the Bochner curvature tensor and apply Lemma 3.1. The Bochner curvature tensor for a Kähler-Einstein metric g is

$$B_{ijk\ell} = R_{ijk\ell} - \frac{s_g}{n(n+2)} \{ (g_{ik}g_{j\ell} - g_{i\ell}g_{jk}) + (J_{ik}J_{j\ell} - J_{jk}J_{i\ell} + 2J_{ij}J_{k\ell}) \}$$

(see for example [6] and [16]) so that a complex space form metric is a Kähler metric whose Bochner curvature tensor vanishes. So, similarly to the proof of Theorem 1 we have

$$(\hat{R}(h_2), h_2)_g = (\hat{B}(h_2), h_2)_g - 2\frac{s_g}{n(n+2)}|h_2|_g^2$$

and the quadratic form Q is then

$$Q_x(h_2, h_2) = -\{\frac{s_g}{n}|h_2|^2 + (\hat{B}(h_2), h_2) - 2\frac{s_g}{n(n+2)}|h_2|^2\}$$
$$= -\{\frac{s_g}{n+2}|h_2|^2 + (\hat{B}(h_2), h_2)\}.$$

Therefore, from the assumption of the theorem there exists $\lambda > 0$ such that $Q_x(h_2, h_2) \ge \lambda |h_2|^2$ for all h_2 . By combining these two arguments we obtain the theorem.

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5. Stability and deformations of complex structures

Let (M, g, J) be a compact Kähler-Einstein manifold and of $s_g < 0$. For the proof of Theorem 3 it suffices to show that there is a $\lambda > 0$ such that

$$\int_M ((D^*D - 2\hat{R})(h), h)_g \, dv_g \ge \lambda \int_M |h|_g^2 dv_g$$

for all h satisfying $\operatorname{tr}_{q} h = 0$ and $\delta_{q} h = 0$.

As discussed in the above we divide the argument into the J-invariant case and the anti-J-invariant case. For the invariant case we have

$$\int_M ((D^*D - 2\hat{R})(h_1), h_1)_g \, dv_g \geq -2rac{s_g}{n} \int_M |h_1|^2 dv_g.$$

On the other hand, for h_2 we set $I \in \Gamma(M; \text{End } (TM))$ by $g(X, IY) = h_2(X, JY)$. Then I is symmetric, i.e., g(IX, Y) = g(X, IY) and anti-commutes with J, i.e., IJ + JI = 0. Therefore we may regard I as a $T^{1,0}M$ -valued (0, 1)-form. By making use of the Weitzenböck formula

$$(D^*D - 2\hat{R})(h_2)(X, JY) = g(X, (\Delta''I)(Y)),$$

where $\Delta'' = \overline{\partial}\overline{\partial}^* + \overline{\partial}^*\overline{\partial}$ is the complex Laplacian associated with the Dolbeault complex of the holomorphic tangent bundle $T^{1,0}M$. See for this, formula 12.93', [1]. Therefore

$$(D^*D - 2\hat{R})(h_2)(X, Y) = -g(X, (\Delta''I)JY)$$

which is equal to $g(X, J(\Delta''I)Y)$, since IJ + JI = 0. Then

$$((D^*D - 2\hat{R})(h_2), h_2) = \sum_{i,j} g(e_i, J(\Delta''I)(e_j))h_2(e_i, e_j) = (\Delta''I, I)_g,$$

since $h_2(e_i, e_j) = g(e_i, JIe_j)$. Therefore

$$\begin{split} \int_{M} ((D^*D - 2\hat{R})(h_2), h_2)_g dv_g &= \int_{M} (\Delta''I, I)_g dv_g \\ &\geq \lambda \int_{M} |I|^2 dv_g = \lambda \int_{M} |h_2|^2 dv_g, \end{split}$$

because the operator Δ'' is positive definite, self-adjoint elliptic and the assumption Ker $\Delta'' = H^1(M; \Theta) = 0$ ensures that the first eigenvalue of Δ'' is positive. Theorem 3 follows from these arguments.

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