

# POSITIVELY RECURRENT MARKOV CHAINS AND THE STEPPING STONE MODEL AS A FLEMING-VIOT PROCESS

By

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**Abstract.** The stepping stone model with infinitely many alleles has been studied in the framework of Fleming-Viot processes by Handa [5]. In this paper, it is investigated the strong-migration limit of the average number of distinct elements in a sample of finite particles in the stationary state of the model, where the results in [5] are applied. To obtain our result, we investigate a problem on positively recurrent Markov chains, and make use of the results in Shiga, Shimizu and Soshi [12].

## 1. Introduction

We will discuss the stepping stone model as a Fleming-Viot process, and the related Markov chains determined by the migration rates. Fleming-Viot processes have been investigated by many authors. See Ethier and Griffiths [1], Ethier and Kurtz [2] and [3]. One of important applications of Fleming-Viot processes is to formulate mathematically the infinite allele model in population genetics which was given by Kimura and Crow [7]. The Fleming-Viot process describing the stepping stone model with infinitely many alleles has been studied by Handa [5]

First, we explain the Fleming-Viot process describing the stepping stone model, and our problem on the measure-valued diffusion. Let  $E$  be the 1-dimensional interval  $[0, 1]$ , and  $S$  be a countable (or finite) set. The set  $P(E)$  of probability measures on  $E$  is endowed with the weak topology. Define  $\mathcal{P}$

$$\mathcal{P} = P(E)^S = \{\tilde{\mu} = \{\mu_k\}_{k \in S}; \mu_k \in P(E), k \in S\},$$

with the product topology. In the following, for a topological space  $X$ ,  $B(X)$  stands for the space of bounded Borel measurable functions on  $X$ . For each positive integer  $k$ ,  $X^k$  denotes the  $k$ -fold product of  $X$ .

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Let  $N$  be the set of positive integers, and  $\theta$  be a positive number. For each  $n \in N$  and any  $f \in B(E^n)$ , define  $L^{(n)}f$  by

$$(L^{(n)}f)(x_1, x_2, \dots, x_n) = \sum_{i=1}^n (L_i f)(x_1, x_2, \dots, x_n),$$

where

$$(L_i f)(x_1, x_2, \dots, x_n) = \frac{\theta}{2} \left( \int_0^1 f(x_1, \dots, x_{i-1}, u, x_{i+1}, \dots, x_n) du - f(x_1, \dots, x_n) \right).$$

For  $n \in N$ ,  $f \in B(E^n)$  and  $1 \leq i < j \leq n$ ,  $\Phi_{ij}^{(n)}f$  is defined by

$$(\Phi_{ij}^{(n)}f)(x_1, x_2, \dots, x_{n-1}) = f(x_1, \dots, x_{j-1}, x_i, x_j, x_{j+1}, \dots, x_{n-1}).$$

Here, in the right-hand side of the above equality, the  $j$ -th variable  $x_j$  is replaced by  $x_i$ , and the variables  $x_{j+1}, x_{j+2}, \dots, x_n$  are replaced by  $x_j, x_{j+1}, \dots, x_{n-1}$  respectively.

Let us introduce a matrix  $\{m_{k'k}\}_{k', k \in S}$ , which describes migration rates. Suppose

$$m_{k'k} \geq 0 \quad \text{if } k \neq k' \quad \text{and} \quad m_{kk} = - \sum_{k': k' \neq k} m_{k'k}.$$

Furthermore, we assume that  $\sup_{k \in S} |m_{kk}| < +\infty$ .

For  $k = (k_1, k_2, \dots, k_n) \in S^n$ ,  $\beta_j k$  and  $\gamma_i(k')k$  are defined by

$$\beta_j k = (k_1, k_2, \dots, k_{j-1}, k_{j+1}, \dots, k_n) \in S^{n-1},$$

and

$$\gamma_i(k')k = (k_1, \dots, k_{i-1}, k', k_{i+1}, \dots, k_n) \in S^n.$$

Namely,  $\beta_j$  eliminates the  $j$ -th coordinate of  $k$ , and  $\gamma_i(k')$  replaces the  $i$ -th element of  $k$  by  $k'$ . For  $k = (k_1, k_2, \dots, k_n) \in S^n$ ,  $\bar{\mu}_k$  denotes  $\mu_{k_1} \times \mu_{k_2} \times \dots \times \mu_{k_n}$  (direct product), and  $\phi_{f,k}(\bar{\mu})$  is defined by  $\phi_{f,k}(\bar{\mu}) = \langle f, \bar{\mu}_k \rangle$  for  $f \in B(E^n)$ .

Now, we are in a position to define the operator  $\mathcal{L}$  by

$$\begin{aligned} (\mathcal{L}\phi_{f,k})(\bar{\mu}) &= \sum_{1 \leq i < j \leq n, k_i = k_j} (\langle \Phi_{ij}^{(n)}f, \bar{\mu}_{\beta_j k} \rangle - \langle f, \bar{\mu}_k \rangle) \\ &\quad + \langle L^{(n)}f, \bar{\mu}_k \rangle + \sum_{i=1}^n \sum_{k' \in S} m_{k'k_i} \langle f, \bar{\mu}_{\gamma_i(k')k} \rangle. \end{aligned} \quad (1.1)$$

It is shown in Handa [5] that the martingale problem for the operator  $\mathcal{L}$  is well posed, and that the  $\mathcal{P}$ -valued diffusion  $\{\mu(t)\}$  determined by  $\mathcal{L}$  has a unique stationary distribution, which is denoted by  $\bar{Q}(d\bar{\mu})$ .

Define

$$I_n = \{\alpha = \{\alpha_k\}_{k \in S}; \alpha_k \in \mathbb{Z}_+, \sum_{k \in S} \alpha_k = n\}, \text{ and } |\alpha| = \sum_{k \in S} \alpha_k.$$

where  $\mathbb{Z}_+$  denotes the set of non-negative integers, and set

$$I = \bigcup_{n=1}^{+\infty} I_n.$$

Consider the mapping  $\Psi : \bigcup_{n=1}^{+\infty} S^n \rightarrow I$ , defined by

$$\Psi(k) = \{\alpha_k\}_{k \in S}, k \in S^n,$$

where  $\alpha_k$  = the cardinality of the set  $\{i; k_i = k\}$ . Let  $\bar{\mu}_\alpha = \prod_{k \in S, \alpha_k > 0} \mu_k^{\alpha_k}$ , where  $\mu_k^{\alpha_k}$  denotes the  $\alpha_k$ -fold product of  $\mu_k$ . If a function  $f \in B(E^n)$  is symmetric, and  $\Psi(k) = \Psi(k') = \alpha \in I_n$ , then  $\langle f, \bar{\mu}_\alpha \rangle$  is well-defined and we have

$$\langle f, \bar{\mu}_k \rangle = \langle f, \bar{\mu}_{k'} \rangle = \langle f, \bar{\mu}_\alpha \rangle.$$

Note that for a symmetric function  $f \in B(E^n)$ ,  $\Phi_{ij}^{(n)} f$  and  $L^{(n)} f$  are symmetric in  $B(E^{n-1})$  and  $B(E^n)$  respectively. In the following, we assume that the migration rates  $m_{k'k}$ ,  $k', k \in S$  are written in the form

$$m_{k'k} = mr_{k'k}, \quad k', k \in S,$$

with a positive constant  $m$ . The stationary distribution  $\bar{Q}(d\bar{\mu})$  is denoted by  $\bar{Q}_m(d\bar{\mu})$ , since it depends on  $m$ . Define  $f_n \in B(E^n)$  by  $f_n(x_1, x_2, \dots, x_n) =$  the number of distinct elements of the sequence  $(x_1, x_2, \dots, x_n)$ . Note that the function  $f_n$  is symmetric.

Our problem on the Fleming-Viot process is to investigate the asymptotic behavior of the integral

$$\mathcal{A}_m(\alpha) := \int \mathcal{F}_\alpha(\bar{\mu}) \bar{Q}_m(d\bar{\mu}), \quad (1.2)$$

as  $m \rightarrow \infty$ , where

$$\mathcal{F}_\alpha(\bar{\mu}) = \langle f_n, \bar{\mu}_\alpha \rangle, \alpha \in I_n.$$

Here, we explain what the quantity  $\mathcal{A}_m(\alpha)$  means. Choose  $n$  genes from the colonies according to  $\alpha \in I_n$ . Namely,  $\alpha_k$  genes are chosen from colony  $k$ , for each  $k \in S$ .  $\mathcal{F}_\alpha(\bar{\mu})$  means the expectation of the number of distinct elements in a sample of  $n$  genes with respect to the measure  $\bar{\mu} \in \mathcal{P}$ . Since  $\mathcal{A}_m(\alpha)$  is the

average of  $\mathcal{F}_\alpha(\bar{\mu})$  with respect to  $\bar{Q}_m(d\bar{\mu})$ ,  $\mathcal{A}_m(\alpha)$ ,  $\alpha \in I_n$ , should be called the average expectation of the number of alleles in a sample of  $n$  genes.

The asymptotic behavior of a quantity for stepping stone models, as  $m \rightarrow \infty$ , is called that of the strong-migration limit. The strong-migration limit in geographically structured population was investigated in Nagylaki [8], Notohara [9] and [10]. Notohara [10] discussed the same problem as ours. In [10], it is assumed that the migration process is a random walk on the  $d$ -dimensional torus  $S$  with lattice points. Under this assumption, he got fruitful results on the strong-migration limit of the stepping stone model. In this paper, the set  $S$  is assumed to be a countable set, so that it will be assumed the finiteness of the moments of the first returning time for the migration process to investigate the asymptotic behavior of  $\mathcal{A}_m(\alpha)$  in the strong-migration limit.

Let us introduce three kinds of Markov chains.

Define  $q_{kk'}$ ,  $k, k' \in S$  by

$$q_{kk'} = r_{k'k}.$$

Let  $q_k = -q_{kk}$ . The minimal Markov chain generated by  $\{q_{kk'}\}_{k,k' \in S}$  is denoted by  $\{x(t), P_k\}$ . It is conservative since  $\{q_k\}_{k \in S}$  is bounded by the assumption on the migration rates. In the following, the Markov chain  $\{x(t), P_k\}$  is sometimes called the 1-particle system. We consider the  $n$ -fold direct product of the 1-particle system  $\{x(t), P_k\}$ , which may be called the  $n$ -particle system. The  $n$ -particle system, whose state space is  $S^n$ , is denoted by  $\{x(t)\}_{t \geq 0}$ . Define the matrix  $\bar{q}_{\alpha\beta}$  on  $I_n$ , by

$$\begin{aligned} \bar{q}_{\alpha\beta} &= \alpha_k r_{k'k} & \text{if } \beta &= \alpha - \epsilon_k + \epsilon_{k'}, \\ &= \sum_k \alpha_k r_{kk} & \text{if } \beta &= \alpha, \\ &= 0 & \text{otherwise,} \end{aligned} \tag{1.3}$$

where  $\epsilon_k = \{\delta_{k,l}\}_{l \in S} \in I_1$ . The Markov chain with state space  $I_n$ , generated by  $\{\bar{q}_{\alpha\beta}\}_{\alpha, \beta \in I_n}$ , is denoted by  $\{\bar{\alpha}(t)\}_{t \geq 0}$ . We see that

$$\{\Psi(x(t))\}_{t \geq 0} \stackrel{\text{law}}{=} \{\bar{\alpha}(t)\}_{t \geq 0}.$$

In this paper, we assume the following.

**ASSUMPTION 1.** The Markov chain  $\{x(t), P_k\}$ , the 1-particle system, is irreducible and positively recurrent, whose stationary distribution is denoted by  $\nu$ .

Under this assumption, the  $n$ -particle system  $\{x(t)\}_{t \geq 0}$  and the Markov chain  $\{\alpha(t)\}_{t \geq 0}$ , mapped by  $\Psi$  from the  $n$ -particle system are both positively recurrent,

whose stationary distributions are  $\nu_n(k) = \prod_{i=1}^n \nu(k_i)$ , for  $k = (k_1, k_2, \dots, k_n) \in S^n$ , and  $\bar{\nu}_n = \frac{n!}{\prod_k \alpha_k!} \nu(k)^{\alpha_k}$  for  $\alpha = (\alpha_k)_{k \in S}$  respectively.

Second, we explain our problem on the 1-particle system  $\{x(t), P_k\}$ . Let us consider the equation on a function  $u_m$

$$m \sum_{k'} q_{kk'} u_m(k') - V(k) u_m(k) = -f(k), \quad k \in S, \quad (1.4)$$

with a positive constant  $m$  and  $f \in B(S)$ , where  $V(k)$  satisfies  $c_1 \leq V(k) \leq c_2$  for some positive constants  $c_1, c_2$ . Recall that the equation (1.4) has a unique bounded solution  $u_m$  if and only if the Markov chain  $\{x(t), P_k\}$  is conservative (Feller [4], Shiga [11]), and it is easy to see that

$$\lim_{m \rightarrow \infty} u_m(k) = \langle f \rangle / \langle V \rangle,$$

where  $\langle g \rangle = \sum_k g(k) \nu(k)$  for  $g \in B(S)$ . Our problem is the following. Under what condition does the sequence  $m(u_m(k) - \langle f \rangle / \langle V \rangle)$  have a finite limit as  $m$  tends to infinity? Our result on the problem mentioned above of the Markov chain  $\{x(t), P_k\}$  is formulated as the next theorem. The expectation with respect to  $P_k$  is denoted by  $E_k[\cdot]$ . Define  $\mathcal{V}$ , a class of bounded functions on  $S$ , by

$$\mathcal{V} = \{V \in B(S); \text{ there exist positive constants } c_1 \text{ and } c_2 \text{ such that } c_1 \leq V(k) \leq c_2, \text{ for any } k \in S\}.$$

Let  $T$  be the first returning time, namely

$$T = \inf\{t > 0; x(t) = x(0) \text{ and } x(s) \neq x(0) \text{ for some } s < t\}.$$

**THEOREM 1.** *There exists a constant  $c$  depending on  $k$ ,  $V$ , and  $f$  for any  $k \in S$ ,  $V \in \mathcal{V}$  and  $f \in B(S)$  such that*

$$\lim_{m \rightarrow \infty} m(u_m(k) - \langle f \rangle / \langle V \rangle) = c,$$

*if and only if*

$$E_{k_0}[T^2] < \infty \quad \text{for some } k_0 \in S. \quad (1.5)$$

*Remark.* Note that the condition  $E_{k_0}[T^2] < \infty$  for some  $k_0 \in S$  implies  $E_k[T^2] < \infty$  for all  $k \in S$ . Since the matrix  $\{q_{kk'}\}$  is defined at the beginning by the migration rates  $\{m_{k'k}\}$  of the stepping stone model, it is assumed in Theorem 1 that  $q_k$  is bounded in  $k \in S$ . However, it suffices to assume on  $\{q_{kk'}\}$  that the minimal Markov chain generated by  $\{q_{kk'}\}$  is conservative, in order to prove Theorem 1 disregarding the stepping stone model.

Here, we apply this result to the problem of the Fleming-Viot process describing a stepping stone model. It is easy to see by using (3.3) and (3.4) below that, under Assumption 1, the equality

$$\lim_{m \rightarrow +\infty} \mathcal{A}_m(\alpha) = 1 + \theta \sum_{l=2}^n \frac{l}{\bar{\nu}_l + (\theta - 1)l} = 1 + \theta \sum_{l=2}^n \frac{1}{(l-1) \sum_k \nu(k)^2 + \theta}$$

holds for  $\alpha \in I_n$ ,  $n \geq 2$ , where  $\bar{\nu}_l = \sum_k \sum_{\alpha \in I_l} \alpha_k^2 \bar{\nu}_l(\alpha) = (l^2 - l) \sum_k \nu(k)^2 + l$ . The right-hand side of the above equality is denoted by  $\mathcal{A}_\infty(\alpha)$ . Recall that the average number of alleles in randomly chosen  $n$  genes for the ordinary infinite allele model is known to be equal to

$$1 + \theta \sum_{l=2}^n \frac{1}{l-1+\theta}.$$

See Hartl and Clark [6]. By Shiga, Shimizu and Soshi [12], we see that the finiteness of the expectation of  $T^2$  for the 1-particle system implies that of the second moment of the first returning time for the  $n$ -particle system. Therefore, under the assumption that  $E_k[T^2] < \infty$  for some  $k \in S$ , the first returning time for the Markov chain  $\bar{\alpha}(t)$ , mapped by  $\Psi$  from the  $n$ -particle system has the finite second moment. Making use of Theorem 1 and the results of Shiga, Shimizu and Soshi [12], we can obtain the next theorem.

**THEOREM 2.** *Assume for the 1-particle system  $\{x(t), P_k\}$  that  $E_k[T^2] < \infty$  for some  $k$ . Then, there exists a constant  $c'$  depending on  $\alpha$  for each  $\alpha \in I_2$  such that*

$$\lim_{m \rightarrow +\infty} m\{\mathcal{A}_m(\alpha) - \mathcal{A}_\infty(\alpha)\} = c'.$$

## 2. Proof of Theorem 1

By the Feynman-Kac formula, we see that the solution  $u_m(k)$  of (1.4) can be written as follows (Shiga [11]):

$$u_m(k) = \frac{1}{m} \int_0^{+\infty} E_k \left[ \exp \left\{ -\frac{1}{m} \int_0^s V(x(u)) du \right\} f(x(s)) \right] ds \quad (2.1)$$

$$= \int_0^{+\infty} E_k \left[ \exp \left\{ -\frac{1}{m} \int_0^{mt} V(x(u)) du \right\} f(x(mt)) \right] dt. \quad (2.2)$$

The ergodic property of the 1-particle system and (2.2) imply that

$$\lim_{m \rightarrow \infty} u_m(k) = \langle f \rangle / \langle V \rangle.$$

Making use of the strong Markov property of the 1-particle system and (2.1), we obtain

$$u_m(k) = \frac{E_k \left[ \int_0^T \exp \left\{ -\frac{1}{m} \int_0^t V(x(u)) du \right\} f(x(t)) dt \right]}{m \left( 1 - E_k \left[ \exp \left\{ -\frac{1}{m} \int_0^T V(x(u)) du \right\} \right] \right)}. \quad (2.3)$$

Define  $\phi(z)$ ,  $\psi(z)$  and  $v(z)$  by

$$\begin{aligned} \phi(z) &= E_k \left[ \int_0^T \exp \left\{ -z \int_0^t V(x(u)) du \right\} f(x(t)) dt \right], \\ \psi(z) &= \frac{1}{z} \left\{ 1 - E_k \left[ \exp \left\{ -z \int_0^T V(x(u)) du \right\} \right] \right\} \equiv \frac{1}{z} \{ 1 - \eta(z) \}, \end{aligned}$$

and

$$v(z) = \frac{\phi(z)}{\psi(z)}$$

for a positive number  $z$ , respectively. Note that  $\phi(0+) = E_k[T]\langle f \rangle$ ,  $\psi(0+) = E_k[T]\langle V \rangle$ , and that  $v(0) = \langle f \rangle / \langle V \rangle$ . First, we assume that  $E_{k_0}[T^2] < \infty$  for some  $k_0 \in S$ . In order to show that the condition (1.5) is sufficient, it suffices to prove that  $v(z)$  is right-differentiable at  $z = 0$ . Since  $E_k[T^2] < +\infty$  for any  $k \in S$ , as mentioned in Remark of Section 1, we see that

$$\lim_{z \downarrow 0} \phi'(z) = -E_k \left[ \int_0^T \int_0^t V(x(u)) du f(x(t)) dt \right].$$

By a simple calculation, we see

$$\psi'(z) = \frac{\eta(z) - 1 - z\eta'(z)}{z^2}.$$

Noting that  $E_k[T^2] < +\infty$  implies the existence of  $\eta''(0+)$ , we get

$$\lim_{z \downarrow 0} \psi'(z) = -\frac{1}{2}\eta''(0+) = -\frac{1}{2}E_k \left[ \left( \int_0^T V(x(u)) du \right)^2 \right].$$

Hence, we conclude that  $v'(0+)$  exists, namely,

$$v'(0+) = \frac{\phi'(0+)\psi(0+) - \phi(0+)\psi'(0+)}{\psi(0+)^2}.$$

To complete the proof of Theorem 1, it is sufficient to show the next lemma.

**LEMMA 2.1.** *Let  $f(k) \equiv 1$ ,  $V(k) = I_{\{k_0\}}(k) + 1$ . Then, the statement that*

$$u_m(k_0) - \frac{1}{1 + \nu(k_0)} = O(1/m), \quad m \rightarrow +\infty,$$

*holds if and only if  $E_{k_0}[T^2] < +\infty$ .*

*Proof.* Noting that  $f(k) \equiv 1$ , we have for the initial state  $k_0$

$$\phi(z) = E_{k_0} \left[ \int_0^T \exp \left\{ -z \int_0^t V(x(u)) du \right\} dt \right].$$

Let  $\zeta$  be the first jumping time, and  $T_0$  be the first passage time to the state  $k_0$ , then

$$\begin{aligned} \phi(z) &= E_{k_0} \left[ \int_0^\zeta \exp \left\{ -z \int_0^t V(x(u)) du \right\} dt \right] + E_{k_0} \left[ \int_\zeta^T \exp \left\{ -z \int_0^t V(x(u)) du \right\} dt \right] \\ &= E_{k_0} \left[ \int_0^\zeta \exp \{-2zt\} dt \right] + E_{k_0} \left[ \int_\zeta^T \exp \{-z(2\zeta + t - \zeta)\} dt \right] \\ &= E_{k_0} \left[ \frac{1}{2z} (1 - e^{-2z\zeta}) \right] + E_{k_0} \left[ e^{-2z\zeta} E_{x_\zeta} \left[ \int_0^{T_0} \exp(-zu) du \right] \right] \\ &= \frac{1}{q_{k_0} + 2z} + \frac{q_{k_0}}{z(q_{k_0} + 2z)} (1 - E_{k_0} [E_{x_\zeta} [e^{-zT_0}]]) \\ &= \frac{1}{z(q_{k_0} + 2z)} (z + q_{k_0} - q_{k_0} E_{k_0} [E_{x_\zeta} [e^{-zT_0}]]) . \end{aligned}$$

Similarly, we see for the initial state  $k_0$  that

$$\psi(z) = \frac{1}{z(q_{k_0} + 2z)} (2z + q_{k_0} - q_{k_0} E_{k_0} [E_{x_\zeta} [e^{-zT_0}]]) .$$

Hence, we obtain

$$v(z) = 1 - \frac{1}{q_{k_0} + 2z} \frac{1}{\psi(z)} ,$$

so that  $v(z)$  is right-differentiable at  $z = 0$  if and only if  $\frac{\psi(z) - \psi(0+)}{z}$  has a finite limit as  $z \downarrow 0$ . Since  $\psi(0+) = E_{k_0} \left[ \int_0^T V(x(u)) du \right]$ , we see that

$$\begin{aligned} \frac{\psi(z) - \psi(0+)}{z} &= \frac{1}{z^2} \left( 1 - E_{k_0} \left[ z \int_0^T V(x(u)) du \right] - E_{k_0} \left[ \exp \left\{ -z \int_0^T V(x(u)) du \right\} \right] \right) \\ &= (-1) E_{k_0} \left[ \int_0^{\int_0^T V(x(u)) du} dt \int_0^t e^{-zs} ds \right] . \end{aligned}$$

Consequently,

$$\lim_{z \downarrow 0} \frac{\psi(z) - \psi(0+)}{z} = (-1) \frac{1}{2} E_{k_0} \left[ \left( \int_0^T V(x(u)) du \right)^2 \right] .$$

Noting that

$$T^2 \leq \left( \int_0^T V(x(u)) du \right)^2 \leq 4T^2 ,$$

we can conclude that there exists a finite limit  $\lim_{z \downarrow 0} \frac{\psi(z) - \psi(0+)}{z}$  if and only if  $E_{k_0} [T^2] < +\infty$ . Thus, the proof of Lemma 2.1 is complete.



### 3. Proof of Theorem 2

First, let us introduce another infinitesimal matrix  $q_{\alpha\beta}^m$  on  $I$  defined by

$$\begin{aligned} q_{\alpha\beta}^m &= \binom{\alpha_k}{2} + \frac{\theta}{2}\alpha_k && \text{if } \beta = \alpha - \epsilon_k \text{ and } |\alpha| \geq 2 \\ &= \alpha_k m r_{k'k} && \text{if } \beta = \alpha - \epsilon_k + \epsilon_{k'}, k \neq k' \\ &= - \sum_{k \in S} \left( \binom{\alpha_k}{2} + \frac{\theta}{2}\alpha_k \right) I_{\{|\alpha| \geq 2\}}(\alpha) + \sum_{k \in S} \alpha_k m r_{kk} && \text{if } \beta = \alpha \\ &= 0 && \text{otherwise,} \end{aligned}$$

where  $\alpha + \beta$  and  $\alpha - \beta$  are defined componentwisely for  $\alpha, \beta \in I$ . Here,  $\binom{n}{k}$  means 0 if  $n < k$ .

Note that the equalities

$$\begin{aligned} (L_i f_n)(x_1, x_2, \dots, x_n) \\ = \frac{\theta}{2} (1 + f_{n-1}(x_1, \dots, x_{i-1}, x_{i+1}, \dots, x_n) - f_n(x_1, x_2, \dots, x_n)), \end{aligned}$$

and

$$\begin{aligned} (\mathcal{L}\mathcal{F}_\alpha)(\bar{\mu}) &= \sum_{k \in S} \binom{\alpha_k}{2} (\langle f_{n-1}, \bar{\mu}_{\alpha - \epsilon_k} \rangle - \langle f_n, \bar{\mu}_\alpha \rangle) + \frac{\theta}{2} |\alpha| \\ &\quad + \sum_{k \in S} \frac{\theta}{2} (\langle f_{n-1}, \bar{\mu}_{\alpha - \epsilon_k} \rangle - \langle f_n, \bar{\mu}_\alpha \rangle) \\ &\quad + \sum_{k \in S} \sum_{k' \in S} \alpha_k m_{k'k} \langle f_n, \bar{\mu}_{\alpha - \epsilon_k + \epsilon_{k'}} \rangle \quad \text{for } |\alpha| = n \geq 2. \end{aligned}$$

hold. Since

$$\int (\mathcal{L}\mathcal{F}_\alpha)(\bar{\mu}) \bar{Q}_m(d\bar{\mu}) = 0,$$

we obtain

$$\sum_{\beta} q_{\alpha\beta}^m \mathcal{A}_m(\beta) + \frac{\theta}{2} |\alpha| = 0 \quad \text{if } |\alpha| \geq 2, \quad (3.1)$$

$$\mathcal{A}_m(\alpha) = 1 \quad \text{if } |\alpha| = 1. \quad (3.2)$$

Let  $(\alpha(t), P_\alpha^m)$ ,  $t \geq 0$ ,  $\alpha \in I$ , be the Markov chain generated by  $\{q_{\alpha\beta}^m\}_{\alpha, \beta \in I}$ .  $E_\alpha^m$  denotes the expectation with respect to  $P_\alpha^m$ . Define  $T$  and  $T_1$  by

$$T = \inf\{t \geq 0; |\alpha(t)| = 1\},$$

and

$$T_1 = \inf\{t \geq 0; |\alpha(t)| < |\alpha(0)|\},$$

where  $\inf \emptyset = +\infty$ . By (3.1) and (3.2), we obtain for  $\alpha \in I_n$ ,  $n \geq 2$ ,

$$\begin{aligned} \mathcal{A}_m(\alpha) &= 1 + E_\alpha^m \left[ \int_0^T \frac{\theta}{2} |\alpha(s)| ds \right] \\ &= \frac{n\theta}{2} E_\alpha^m [T_1] + E_\alpha^m [\mathcal{A}_m(\alpha(T_1))]. \end{aligned} \quad (3.3)$$

Observe that  $\mathcal{A}_m^1(\alpha) := E_\alpha^m [T_1]$  satisfies

$$\sum_{\beta \in I_n} \bar{q}_{\alpha\beta} \mathcal{A}_m^1(\beta) - \frac{1}{m} \frac{1}{2} \left( \sum_i \alpha_i^2 + (\theta - 1)n \right) \mathcal{A}_m^1(\alpha) = -\frac{1}{m}, \quad (3.4)$$

with the matrix  $\{\bar{q}_{\alpha\beta}\}_{\alpha, \beta \in I_n}$  defined by (1.3). Since the minimal Markov chain  $\{\bar{\alpha}(t), P_\alpha\}$  on  $I_n$  generated by the matrix  $\{\bar{q}_{\alpha\beta}\}_{\alpha, \beta \in I_n}$  is conservative and positively recurrent, as explained in Section 1. Therefore, we can apply Theorem 1 to  $\{\bar{\alpha}(t), P_\alpha\}$ . Let  $T$  be the first returning time of  $\{\bar{\alpha}(t), P_\alpha\}$ , then by the results of Shiga, Shimizu and Soshi [12], as mentioned in Section 1, we see that  $E_\alpha[T^2] < +\infty$ , where  $E_\alpha[\cdot]$  denotes the expectation with respect to  $P_\alpha$ . Thus, we obtain the next lemma.

**LEMMA 3.1.** *For each  $\alpha \in I_n$ , there exists a finite limit  $\lim_{m \rightarrow \infty} m(\mathcal{A}_m^1(\alpha) - \frac{2}{\bar{v}_n + (\theta - 1)n})$ .*

If  $\alpha \in I_2$ , the second term of the right-hand side of (3.3) is equal to 1, so that Lemma 3.1 proves Theorem 2.

#### 4. An Example

Here, we present a Markov chain for which one can get the critical parameter for the second moment of the first returning time  $T$  to be finite. Furthermore, for the example, it is made clear the speed of convergence of the solution  $u_m$  for the equation (1.4) with specified functions  $V$  and  $f$ , when the second moment of  $T$  is infinite. Proposition 4.1 and Lemma 4.2, given below, are essentially due to Professor T. Shiga.

Consider a continuous time minimal Markov chain  $\{x(t), P_i\}$  in  $S = Z_+$  which is governed by the following infinitesimal matrix  $Q = \{q_{i,j}\}$ .

$$q_{i,j} = \begin{cases} 1 & (\text{if } j = i + 1) \\ \alpha_1 & (\text{if } j = 0) \\ -(1 + \alpha_i) & (\text{if } j = i > 0), \end{cases}$$

and

$$q_{0,j} = \delta_{1,j} \quad (j \geq 1) \quad q_{0,0} = -1,$$

$$q_{i,j} = 0 \quad \text{for any other } i, j \in S.$$

Recall that  $T$  denotes the first returning time of the continuous time Markov chain  $\{x(t), P_i\}$ . Then it is easy to show that

$$E_0(e^{-\lambda T}) = \sum_{i=1}^{\infty} \frac{\alpha_i}{1+\lambda} \prod_{k=1}^i \frac{1}{1+\lambda+\alpha_k}, \quad (4.1)$$

from which it follows that  $\{x(t), P_i\}$  is recurrent if and only if

$$\sum_{i=1}^{\infty} \alpha_i = \infty. \quad (4.2)$$

Now we specialize  $\{\alpha_i\}$  as follows:

$$\alpha_i = \frac{\kappa}{i} \quad (i \geq 1), \quad (4.3)$$

with a constant  $\kappa > 0$ . Inserting this to (4.1) we see

$$\begin{aligned} E_0(e^{-\lambda T}) &= \sum_{n=1}^{\infty} \frac{\kappa \Gamma(n) \Gamma(\frac{\kappa}{1+\lambda} + 1)}{(1+\lambda)^{n+1} \Gamma(\frac{\kappa}{1+\lambda} + n + 1)} \\ &= \sum_{n=1}^{\infty} \frac{\kappa}{(1+\lambda)^{n+1}} B\left(n, \frac{\kappa}{1+\lambda} + 1\right), \end{aligned}$$

where  $B(\alpha, \beta)$  denotes the Beta function, so that it holds

$$E_0(e^{-\lambda T}) = \frac{\kappa}{1+\lambda} \int_0^1 \frac{t^{\frac{\kappa}{1+\lambda}}}{\lambda+t} dt. \quad (4.4)$$

**PROPOSITION 4.1.**

$$E_0(T^\alpha) < \infty \quad \text{if and only if} \quad \alpha < \kappa.$$

*Proof.* Using Taylor's expansion of  $1/(\lambda+t)$ , from (4.4) one gets

$$\begin{aligned} E_0(e^{-\lambda T}) &= \kappa \sum_{m=0}^n \frac{(-\lambda)^m}{\kappa - m(\lambda+1)} + \frac{\kappa(-\lambda)^{n+1}}{\lambda+1} \int_0^1 \frac{t^{\frac{\kappa}{1+\lambda} - n - 1}}{\lambda+t} dt \\ &= \kappa \sum_{m=0}^n \frac{(-\lambda)^m}{\kappa - m(\lambda+1)} + \frac{\kappa(-1)^{n+1}}{\lambda+1} \lambda^{\kappa/(\lambda+1)} \int_0^{1/\lambda} \frac{t^{\frac{\kappa}{1+\lambda} - n - 1}}{1+t} dt. \end{aligned}$$

Let  $n < \kappa \leq n+1$  ( $n \in \mathbb{Z}_+$ ). Denoting the last term by  $b(\lambda)$ , we obtain that if  $n < \kappa < n+1$ ,

$$\lim_{\lambda \searrow 0} (-1)^{n+1} \frac{b(\lambda)}{\lambda^\kappa} = \kappa \int_0^\infty \frac{t^{\kappa-n-1}}{1+t} dt, \quad (4.5)$$

and if  $\kappa = n + 1$ ,

$$b(\lambda) \sim \kappa(-\lambda)^{n+1} \log \frac{1}{\lambda} \quad (\lambda \searrow 0). \quad (4.6)$$

Therefore, by applying the following lemma, the proof of Proposition 4.1 is completed.

**LEMMA 4.1.** *For a nonnegative random variable  $Y$ , let*

$$\varphi(\lambda) = E(e^{-\lambda Y}) \quad (\lambda > 0).$$

*Suppose that  $\varphi(\lambda)$  has the form,*

$$\varphi(\lambda) = \sum_{m=0}^n a_m \lambda^m + b(\lambda),$$

*where for a  $\kappa \in (n, n+1]$ , either*

$$b(\lambda) \sim \beta_0 \lambda^\kappa, \quad \text{as } \lambda \downarrow 0,$$

*with  $\beta_0 (\neq 0) \in \mathbb{R}$  or*

$$\lim_{\lambda \searrow 0} (-1)^{n+1} \frac{b(\lambda)}{\lambda^\kappa} = \infty$$

*and for every  $\beta < \kappa$*

$$\lim_{\lambda \searrow 0} (-1)^{n+1} \frac{b(\lambda)}{\lambda^\beta} = 0.$$

*Then it holds that*

$$E(Y^\alpha) < \infty \quad \text{for } \alpha < \kappa, \quad \text{and} \quad E(Y^\kappa) = \infty.$$

Proof of Lemma 4.2 is standard, so it is omitted.

We see by Proposition 4.1 that  $\{x(t), P_i\}$  is positively recurrent if and only if  $\kappa > 1$ , so that we investigate the case of  $\kappa > 1$  in the following. Let us consider (1.4) for this example, with specified functions  $V$  and  $f$ . Let  $V(i) = 1 + I_{\{0\}}(i)$ ,  $f(i) \equiv 1$ , where  $I_{\{0\}}$  is the indicator function on the set  $\{0\}$ . Namely, we consider the equation

$$m \sum_j q_{ij} u_m(j) - V(i) u_m(i) = -1, \quad (4.7)$$

with a positive constant  $m$ . The unique solution  $u_m$  of (4.7) satisfies

$$\lim_{m \rightarrow +\infty} u_m(0) = \frac{1}{1 + \nu(0)},$$

with  $\nu(0) = 1 - \frac{1}{\kappa}$ , and, by the same argument as the proof of Lemma 2.1, it is shown that

$$u_m(0) = 1 - \frac{1}{1+2z} \frac{1}{\psi(z)}, \quad (4.8)$$

where  $z = 1/m$ , and

$$\psi(z) = \frac{1}{z(1+2z)} (2z + 1 - E_1[e^{-zT_0}]), \quad (4.9)$$

where  $T_0$  denotes the first passage time to the state 0. Note that

$$E_1[e^{-zT_0}] = \kappa \int_0^1 t^{\frac{\kappa}{1+z}} \frac{1}{z+t} dt. \quad (4.10)$$

Then, we obtain the next statement.

**PROPOSITION 4.2.** *If  $\kappa > 2$ ,*

$$u_m(0) - \frac{1}{1+\nu(0)} \sim -k_1(\kappa) \frac{1}{m}, \quad m \rightarrow +\infty,$$

*if  $\kappa = 2$ ,*

$$u_m(0) - \frac{1}{1+\nu(0)} \sim -\frac{2}{9} \frac{1}{m} \log m, \quad m \rightarrow +\infty,$$

*and if  $1 < \kappa < 2$ ,*

$$u_m(0) - \frac{1}{1+\nu(0)} \sim -k_2(\kappa) \left(\frac{1}{m}\right)^{\kappa-1}, \quad m \rightarrow +\infty,$$

*where*

$$k_i(\kappa) > 0, \quad (i = 1, 2), \quad \lim_{\kappa \downarrow 2} k_1(\kappa) = +\infty, \quad \lim_{\kappa \uparrow 2} k_2(\kappa) = +\infty,$$

*and  $\nu(0) = 1 - 1/\kappa$ .*

*Proof.* Recall the equality in the proof of Proposition 4.1,

$$E_0[e^{-zT}] = \kappa \sum_{l=0}^n \frac{(-z)^l}{\kappa - l(z+1)} + b(z), \quad (4.11)$$

where

$$b(z) = \frac{\kappa(-1)^{n+1}}{z+1} z^{\frac{\kappa}{z+1}} \int_0^{\frac{1}{z}} \frac{t^{\kappa/(1+z)-n-1}}{1+t} dt. \quad (4.12)$$

We prove Proposition 4.2 only for the case of  $\kappa = 2$ , because one can prove the assertions similarly for  $\kappa > 2$  and  $1 < \kappa < 2$  making use of (4.8), (4.9), (4.10), (4.11) and (4.12). For the case of  $\kappa = 2$ , by (4.11) with  $n = 1$ , we see that

$$a(z) \equiv E_0[e^{-zT}] = 1 - 2z + b(z) + O(z^2), \quad z \downarrow 0,$$

and that

$$b(z) \sim 2z^2 \log \frac{1}{z}, \quad z \downarrow 0.$$

Note that we have by (4.9)

$$\psi(z) = -\frac{1}{1+2z} \frac{a(z)-1}{z} + \frac{2-a(z)}{1+2z},$$

and therefore  $\psi(0+) = 3$ . Consequently, we have

$$\begin{aligned} \frac{\psi(z) - \psi(0+)}{z} &= -\frac{1}{1+2z} \left( \frac{b(z)}{z^2} + O(1) \right), \\ &\sim -2 \log \frac{1}{z}. \end{aligned}$$

Hence, we conclude that

$$\frac{v(z) - v(0+)}{z} \sim -\frac{2}{\psi(0+)^2} \log \frac{1}{z} = -\frac{2}{9} \log \frac{1}{z}, \quad z \downarrow 0.$$

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