

1 Deformation mechanism of capsule-type hydrogen-storage-alloy actuator

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## 12 **Abstract**

13 A capsule-type actuator using hydrogen storage alloy (HSA) was fabricated and its

14 response behavior was investigated by experiment and finite element analysis. The

15 capsule-type HSA actuator consists of an HSA foil placed on the inner wall of a capsule.

16 In this paper, the displacement change of the actuator made of palladium foil and epoxy

17 resin was experimentally measured in hydrogen atmosphere. As a result, it displaced with

18 the decrease in the capsule height when hydrogen gas was charged, and vice versa.

19 Displacement amount and response time of the actuator depended on the hydrogen

20 pressure during the pressurization, while the response time was independent of the

21 pressure during the evacuation. Also, the HSA actuator's deformation with hydrogen

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22 diffusion was simulated using the finite element method. It revealed that hydrogen content  
23 in the HSA foil was decreased by high compressive stress, indicating the hydrogen atom  
24 was absorbed as a solid solution phase in the HSA foil even when a surrounding hydrogen  
25 pressure was higher than the phase transition pressure.

## 26 **Keywords**

27 metal hydride; stress induced diffusion; finite element method; structural-diffusion  
28 coupled analysis; palladium

## 29 **1. Introduction**

30 A hydrogen-storage-alloy (HSA) actuator is an actuator using volume change of  
31 hydrogen storage alloy which reversibly absorbs hydrogen and expands with an increase  
32 in its surrounding H<sub>2</sub> pressure. The actuator is driven by the change in H<sub>2</sub> pressure and  
33 anticipated to be used in a hydrogen society that uses hydrogen as a major source of power  
34 [1, 2].

35 Two types of the HSA actuators have been developed, i.e. a unimorph-type actuator [3-  
36 10] and a capsule-type actuator [11, 12]. In the unimorph actuator, an HSA foil is bonded  
37 with a substrate. Expansion of the HSA foil in H<sub>2</sub> causes bending and twisting motion.  
38 Nishi et al. [3], for example, fabricated a unimorph actuator consisting of a LaNi<sub>5</sub> foil and  
39 a polyimide substrate for medical application. It showed a bending strain of ~0.1% under  
40 hydrogen of 0.12 MPa, and its shape recovered after the evacuation.

41 The capsule-type actuator [11, 12] was proposed for a super multi-link space  
42 manipulator to capture space debris [13]. The manipulator has multiple joints with the  
43 actuators to realize the flexible grasping. Another possible application is a smart regulator  
44 mounted on a fuel-cell vehicle, in which the flux of hydrogen is automatically controlled

45 depending on the pressure. In these structural applications, an actuator is required to have  
46 high stiffness and lightweight in addition to good energy efficiency.

47 The capsule-type actuator has an axisymmetric structure consisting of a hollow capsule  
48 and an HSA foil which is partially placed on the inner wall of the capsule (Figure 1),  
49 which realizes high stiffness and lightweight. The HSA foil expands when H<sub>2</sub> gas is  
50 charged into the capsule, causing the deformation of the capsule. The height change of  
51 the capsule is used as an actuator's displacement. Goto et al. [11] estimated its static  
52 performance by a finite element method (FEM) assuming that the HSA foil isotropically  
53 absorbs hydrogen. They showed the capsule height decreases when the foil absorbs  
54 hydrogen (Figure 1). This concept was experimentally confirmed with a capsule-type  
55 actuator made of palladium and ABS resin [12]. The displacement was, however, much  
56 smaller than that expected from the FEM simulation.

57 The performance of the HSA actuator is strongly affected not only by its shape [7, 11]  
58 but also by hydrogen absorption characteristics of the hydrogen storage alloy used. Honjo  
59 et al. [4] and Nakai et al. [5] investigated deformation of unimorph actuators using La-Ni  
60 and Pd-Ni alloys, respectively, and showed their displacement amounts and response rates  
61 were dependent on the composition of alloys. It is also important to investigate the  
62 swelling ability during hydrogen absorption in a hydrogen storage alloy itself. Some  
63 studies have been performed using powder as a sample, in which various measurement  
64 methods were adopted, for example, optically [14, 15], electrically [16, 17], and with  
65 strain gauges on a sample holder [18-20].

66 A hydrogen absorption behavior of a hydrogen storage alloy is also affected by stress.  
67 Hydrogen content thermodynamically decreases with increasing compressive stress, and  
68 vice versa [21, 22]. The diffusion is promoted not only by a hydrogen content gradient

69 (Fick's law) but also by a stress gradient [23, 24]. Many researchers [25-28] simulated  
70 diffusion of hydrogen atoms in a hydrogen storage alloy under stress, which revealed that  
71 the so-called stress-induced diffusion (SID) plays an important role in the diffusion  
72 behavior. In the actuator, the stress effect cannot be negligible because it receives a  
73 reaction force during operation.

74 In this paper, the response behavior of the capsule-type HSA actuator was investigated.  
75 Displacement of an actuator made of palladium and epoxy resin was measured with  
76 changing H<sub>2</sub> pressure. To reveal its internal state, a numerical simulation was performed  
77 and compared with the experimental results.

## 78 **2. Experiment**

79 The fabricated sample of the capsule-type HSA actuator is shown in Figure 2. It was  
80 almost spherical with outer diameter of  $22.1 \pm 0.4$  mm on the equator (Figure 2a-c). From  
81 a cross-sectional observation after the experiment (Figure 2d), the sample was confirmed  
82 to have a hollow cylindrical structure with inner diameter of  $\sim 13.2$  mm. Its design  
83 dimensions are shown in Figure 2e.

84 It was fabricated as follows (Figure 3). Palladium foil with 0.08 mm in thickness and  
85 epoxy resin were used as the HSA foil and capsule, respectively. The foil (Tanaka  
86 Kikinzoku Kogyo K.K.) of 15 x 42 x 0.08 mm was annealed in the air at 400°C for 2  
87 hours (1 & 2) to remove the residual stress induced by rolling. The annealed foil was  
88 looped on a core made of wax, which is used to make the sample hollow (3). The epoxy  
89 resin was prepared by mixing Epon828 (Polyscience, Inc.) with diethylenetriamine (Alfa  
90 Aesar) (4), followed by vacuum degassing for 5 min (5). It was cast between a silicone  
91 mold and the core with the foil (6). After solidified for 24 hours at room temperature, it

92 was heated at 80°C in the air to remove the core (7). The sample was finally cleaned with  
93 2-propanol. Although the palladium foil formed an oxide layer during the annealing  
94 (Figure 2a), it was cleaned after hydrogen charge, and its color recovered to silver (Figure  
95 2c). No delamination between the HSA foil and the capsule was observed after the  
96 experiments (Figure 2d).

97 The sample was fixed perpendicular to the ground in a pressure vessel. Then, a piston  
98 was put on the sample. Displacement of the piston i.e. height change of the sample  $u$  was  
99 measured with a laser displacement sensor (LK-G85, Keyence) through a viewing  
100 window of the pressure vessel at room temperature. The measurement apparatus is  
101 described in detail in [14]. The purity of hydrogen gas (Taiyo Nippon Sanyo Co.) was  
102 99.999% or more. Before the measurement, hydrogen was introduced into the evacuated  
103 vessel at 5 atm for  $3 \times 10^4$  s, followed by the evacuation, for an activation treatment.

104 The hydrogen pressure was rapidly changed by opening a valve connected to a  
105 hydrogen tank or a rotary vacuum pump at  $t = 0$  s, and then, the vessel was kept at the  
106 constant pressure to measure the response time. It is noted that hydrogen was charged not  
107 only in the capsule but inside the whole of the pressure vessel. The measurement interval  
108 was 0.1 or 1 s during the pressurization and 1 s during the evacuation.

### 109 **3. Experimental results**

110 Figure 4 shows the displacement change in time during the pressurization and the  
111 following evacuation at the introduced hydrogen pressure  $P_{intro}$  of 4.90 atm. The  
112 displacement  $u_i$  at  $t = 0$  s was set to 0.0  $\mu\text{m}$ . The positive value represents an increase in  
113 the capsule height. During the pressurization (Figure 4a), hydrogen pressure yielded 90%  
114 of  $P_{intro}$  within 6 s. At the same time, the displacement increased by 7.5  $\mu\text{m}$  ( $= u_P$ ) in

115 proportion to the pressure. The displacement continued increasing after the pressure  
116 change, and it yielded a maximum value  $u_{max}$  of 10.0  $\mu\text{m}$  at  $t = 245$  s, followed by  
117 decreasing. The final displacement  $u_f$  was  $-6.5 \pm 0.2$   $\mu\text{m}$ , where the error range represents  
118 the standard deviation calculated from the values at last 100 s.

119 The first increase in the displacement at the same time as the pressurization  $u_P$  was  
120 caused by the changes in a refractive index in the vessel and shrinkage of the capsule due  
121 to hydrostatic pressure. The refractive index depends on gas species and its pressure and  
122 temperature. Its contribution is estimated at 1.0  $\mu\text{m}/\text{atm}$  from Snell's law. The  
123 displacement by the hydrostatic pressure was estimated at  $-0.2$   $\mu\text{m}/\text{atm}$  by a finite element  
124 analysis. The total contribution of the refractive index and hydrostatic pressure is,  
125 therefore, 0.8  $\mu\text{m}/\text{atm}$ , which roughly agrees with  $u_P$ . The refractive index change does  
126 not affect the deformation in the practical environment, because hydrogen is charged only  
127 in the capsule. The following increase in displacement is caused by thermal expansion.  
128 The enthalpies of hydrogen absorption in solid solution and the phase transition into  
129 hydride are  $-4620$  and  $-9440$   $\text{cal}/\text{molH}_2$ , respectively in palladium [29], indicating both  
130 reactions are exothermal reactions. Therefore, the temperature in the HSA foil increases  
131 during the pressurization. A net actuator's displacement  $U$  is calculated as  $|u_f - u_P|$ , i.e.,  
132 14.0  $\mu\text{m}$ .

133 During the evacuation (Figure 4b), the displacement almost recovered to 0. It decreased  
134 by 7.5  $\mu\text{m}$  soon after the evacuation ( $u_P$ ). It is due to the change in hydrostatic pressure  
135 and refractive index as well as during the pressurization. Under the constant pressure of  
136 0.00 atm, the displacement gradually increased. The final displacement  $u_f$  was  $1.3 \pm 0.2$   
137  $\mu\text{m}$ . The net displacement  $U$  was calculated as 15.0  $\mu\text{m}$ , which is comparable to that  
138 during the pressurization. The difference of  $U$  between during the pressurization and

139 evacuation is considered to involve the effects of the plastic deformation of the capsule.

140 The similar tendencies were obtained at other  $P_{intro}$ . All data are tabulated in Table 1.

141

142 Table 1 Displacement values at each  $P_{intro}$ .

$P_{intro}$	Pressurization						Evacuation				
	$u_i$	$u_P$	$u_{max}$	$u_f$	$U$	$t_{80}$	$u_i$	$u_P$	$u_f$	$U$	$t_{80}$
[atm]	[ $\mu\text{m}$ ]	[ $\mu\text{m}$ ]	[ $\mu\text{m}$ ]	[ $\mu\text{m}$ ]	[ $\mu\text{m}$ ]	[s]	[ $\mu\text{m}$ ]	[ $\mu\text{m}$ ]	[ $\mu\text{m}$ ]	[ $\mu\text{m}$ ]	[s]
1.16	0.0	0.8	2.9	-4.1	4.9	3348	-4.9	-2.0	-1.0	5.9	2358
1.97	0.0	2.2	4.7	-4.7	6.9	4936	-12.1	-2.8	-6.3	8.6	4045
2.88	0.0	5.4	6.6	-7.5	12.9	1966	-7.9	-4.3	-1.8	10.4	2333
2.91	0.0	5.8	6.6	-4.3	10.1	3055	-5.7	-5.0	-1.7	9.1	4071
4.89	0.0	7.9	10.3	-7.3	15.2	1911	-7.9	-7.1	-1.4	13.7	3370
4.90	0.0	7.5	10.0	-6.5	14.0	1521	-6.2	-7.5	1.3	15.0	2183

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144 Figure 5 shows a variation of the net displacement  $U$  with  $P_{intro}$ .  $U$  during the  
 145 evacuation was comparable to that during the pressurization at all  $P_{intro}$ , indicating that  
 146 the actuator reversibly deformed. Larger net displacement was obtained at higher pressure,  
 147 and its increment decreased with the pressure increase.

148 The dependence of response time on pressure was different during the pressurization  
 149 and evacuation as shown in Figure 6. The 80% response time  $t_{80}$  represents a time to yield  
 150 80% of the net displacement. During the pressurization (Figure 6a),  $t_{80}$  decreased with  
 151  $P_{intro}$ . The response time at  $P_{intro} < 2$  atm was half of that at  $P_{intro} = 5$  atm. During the  
 152 evacuation (Figure 6b),  $t_{80}$  was independent of  $P_{intro}$  although the large scatter was

153 observed. The similar scattering was also observed in our previous measurement using  
154 the same measurement system [14]. The evacuation pressure and purity of the gas in the  
155 pressure chamber depend on the ambient temperature and humidity, leak property of the  
156 chamber and performance of the vacuum pump. The scattering will become small if the  
157 evacuation system is improved.

#### 158 **4. Numerical simulation of HSA actuator's response**

159 The diffusion-mechanical coupling analysis revealed an internal state of the actuator.  
160 In the simulation, hydrogen diffusion was driven not only by concentration gradient but  
161 also by stress gradient. Most of the papers relevant to the stress-induced diffusion [21-28]  
162 focus on hydrogen embrittlement in steel and hydrogen permeable membranes, in which  
163 hydrogen content is small enough to neglect the effect of volume and activity change with  
164 hydrogen content. In the actuator's applications, these effects should be taken into account  
165 since the hydrogen content can be high. The diffusion equation used in this work is as  
166 follows.

$$167 \quad J = -D''\nabla n_H + \frac{D'n_H V_H}{RT} \nabla \sigma_h \quad (1)$$

168 where  $J$ : flux of hydrogen,  $D'$  and  $D''$ : modified hydrogen diffusion coefficients,  $n_H$ :  
169 molar ratio of hydrogen to metal (H/M),  $V_H$ : partial molar volume of hydrogen,  $R$ :  
170 universal gas constant,  $T$ : temperature, and  $\sigma_h$ : hydrostatic stress. The derivation of Eq.  
171 (1) is shown in Appendix A. Goto et al. [28] showed that the stress gradient does not  
172 significantly contribute to the diffusion behavior compared to the activity change with  
173 hydrogen content, while the equilibrium hydrogen content is affected by the stress. A  
174 surface flux of the HSA foil was calculated from a difference between chemical potentials  
175 of hydrogen gas and atoms at the surface. The detailed analytical condition is shown in



176 Appendix B. The displacement amount under 1 atm of hydrogen was 6.9  $\mu\text{m}$  (Figure B.2),  
177 which almost agreed with the experimental result (4.9  $\mu\text{m}$  at 1.16 atm). The difference  
178 between the experiment and simulation seemingly comes from local stress relaxation and  
179 incompleteness of the sample shape. The stress is locally reduced at corners of the HSA  
180 foil due to plastic and viscoelastic deformation in the epoxy resin, which decreases the  
181 displacement. The plasticity and viscoelasticity in the capsule were not taken into account  
182 in the present simulation. Also, the capsule shape of the sample is not completely  
183 spherical. Investigation of these effects on the displacement is the future task.

184 The simulation revealed the internal state of the actuator. Figure 7 shows the time  
185 change of hydrogen content  $n_H$  and compressive hydrostatic stress  $\sigma_h$  in the HSA foil  
186 when a deformation ratio was 1.3, 14.9, 83.4 and 100%, respectively. The deformation  
187 ratio was calculated by dividing displacement by final displacement. The transverse axis  
188 represents the thickness of the foil in the radial direction on the equator position.  $n_H$   
189 gradually increased with time near the surface and diffused into the inside. The foil was  
190 subjected to the higher compressive stress with the increase in the hydrogen content  
191 because its surrounding capsule resisted the deformation.  $n_H$  and  $\sigma_h$  at  $t = 3000$  s were  
192  $5.4 \times 10^{-3}$  and 22.3 MPa, respectively.  $\sigma_h$  at the boundaries was lower than inside because  
193 the stress in the radial direction becomes nearly 0 on the boundaries.

## 194 **5. Discussions**

195 The simulation indicates that hydride is not formed in the HSA foil during the  
196 deformation. Under free stress, the phase transition pressure between solid solution and  
197 hydride in palladium is  $1 \times 10^{-2}$  atm [30]. Therefore, at 1 atm, the hydride is not formed in  
198 palladium, and its equilibrium hydrogen content is  $n_H = 0.72$  [31]. However,  $n_H = 0.0054$

199 in the simulation, which is more than 100 times smaller than 0.72. The palladium-  
200 hydrogen system forms solid-solution state at  $n_H < 0.03$ . This is why the deformation in  
201 the experiment was much less than that expected with the technique used in [11]. The low  
202 equilibrium hydrogen content is due to the high compressive stress.

203 The response time decreased with the increase in the introduced hydrogen pressure  
204  $P_{intro}$  during the pressurization, while it is independent during the evacuation. The same  
205 tendency was observed when hydrogen was charged into palladium powder [14]. The  
206 driving force of hydrogen absorption/discharge increases with the pressure difference  
207 between applied pressure  $P_{ap}$  and equilibrium pressure  $P_{eq}$  [32], where  $P_{ap} = P_{intro}$  during  
208 the pressurization and  $P_{ap} = 0$  atm during the evacuation. The constant response time  
209 independent of the pressure at the time of evacuation implies the existence of the rate-  
210 determining step which gives constant  $P_{eq}$  such as interfacial reaction.

211 It took a few thousand seconds during both of the pressurization and evacuation to  
212 complete the deformation. In contrast, when the volume change of palladium powder was  
213 measured under free stress [14], the response time during the evacuation was longer than  
214 that during the pressurization. The difference comes from the phase of palladium-  
215 hydrogen system: solid solution in the actuator, while hydride under free stress. Ulvestad  
216 et al. [33] pointed out that a hydride on palladium surface decreases a hydrogen discharge  
217 rate. In this study, the simulation indicates the hydride was not formed during the  
218 deformation due to the compressive stress, hence, the response time during the evacuation  
219 was similar during the pressurization.

220 The temperature also affects the response time of the actuator. The reaction rate  
221 increases with temperature. During the pressurization, the response time can be increased  
222 by the exothermal reaction of absorption. The evaluation of the thermal effect is the future

223 task.

## 224 **6. Conclusion**

225 In this paper, a capsule-type hydrogen-storage-alloy (HSA) actuator made of palladium  
226 and epoxy resin as an HSA foil and a capsule was fabricated, and its response behavior  
227 in hydrogen was investigated. Moreover, a mechanical finite element analysis coupled  
228 with diffusion problem was performed to investigate its internal state. The main  
229 conclusions are summarized as follows;

- 230 1. The displacement was observed when hydrogen gas was charged. By the following  
231 evacuation, it recovered to almost 0.
- 232 2. Larger displacement was obtained at higher pressure, and its increment decreased with  
233 the pressure increase.
- 234 3. The displacement rate increased with H<sub>2</sub> pressure during the pressurization, while it  
235 was independent of the pressure during the evacuation.
- 236 4. The solid solution phase is formed inside the palladium foil during deformation due to  
237 compressive stress.

238 Thus, the response behavior of the capsule-type HSA actuator is strongly affected by  
239 the stress, especially in terms of the absorbed hydrogen amount. The performance of the  
240 actuators can be improved by designing to decrease stress in the hydrogen storage alloy.

241 Although the epoxy resin was used for the capsule in this work due to easy forming, it  
242 is not suitable in practical use since hydrogen penetrates the resin. The candidate materials  
243 for the capsule are an aluminum alloy and stainless steel, which has high resistance to  
244 hydrogen penetration. It is a future task to develop the fabrication process of the capsule-  
245 type actuator with these materials. However, the knowledge obtained in this work is

246 useful for the design of all HSA actuators.

## 247 **Appendix A. Derivation of diffusion equation**

248 Chemical potential,  $\mu_H$  including a stress term is represented as Eq. (A.1).

$$249 \quad \mu_H = \mu_H^0 + RT \ln a_H - V_H \sigma_h \quad (\text{A. 1})$$

250 where  $a_H$  is the activity of hydrogen.  $a_H$  is represented with the hydrogen content  $n_H$  and  
251 the activity coefficient  $\gamma_n$  as  $a_H = \gamma_n n_H$ . Hydrogen flux in a hydrogen storage alloy  $J$  is  
252 proportional to a chemical potential gradient.

$$253 \quad J = -\alpha \nabla \mu_H = -\alpha \left\{ \frac{RT}{n_H} \left( 1 + \frac{\partial \ln \gamma_H}{\partial \ln n_H} \right) \nabla n_H - V_H \nabla \sigma_h \right\} \quad (\text{A. 2})$$

254 In case of  $n_H \ll 1$  and free stress, Eq. (A.2) corresponds to Fick's equation.

$$255 \quad J = -\alpha \frac{RT}{n_H} \nabla n_H = -D \nabla C_H \quad (\text{A. 3})$$

$$256 \quad \therefore \alpha = \frac{D n_H}{RT} \frac{dC_H}{dn_H} \quad (\text{A. 4})$$

257 where  $C_H$  is the molar concentration and  $D$  is Fick's diffusion coefficient. A lattice volume  
258 of palladium increases in proportion to  $n_H$  at  $n_H < 0.75$  [34]. The relationship between  $C_H$   
259 and  $n_H$ , therefore, is represented as Eq. (A.5) using  $V_H$  and molar volume of palladium  
260  $V_M$ .

$$261 \quad C_H = \frac{n_H}{V_M + n_H V_H} \quad (\text{A. 5})$$

262 From Eqs. (A.2), (A.4) and (A.5),

$$263 \quad J = -D \frac{V_M}{(V_M + n_H V_H)^2} \left( 1 + \frac{\partial \ln \gamma_H}{\partial \ln n_H} \right) \nabla n_H + \frac{D n_H V_H}{RT} \frac{V_M}{(V_M + n_H V_H)^2} \nabla \sigma_h \quad (\text{A. 6})$$

264 Eq. (1) is derived by defining  $D'$  and  $D''$  as

$$265 \quad D' = D \frac{V_M}{(V_M + n_H V_H)^2}, D'' = D' \left( 1 + \frac{\partial \ln \gamma_H}{\partial \ln n_H} \right) \quad (\text{A. 7})$$

266  $D'$  includes the effect of the volume change during absorption, and  $D''$  reflects the change

267 in the activity in addition to the volume change.

268 The activity coefficient was determined from a PCT diagram of the palladium-  
 269 hydrogen system [31]. The chemical potential of molecules in the gas  $\mu_{H_2}$  is same as that  
 270 of atoms in the alloy in the equilibrium condition, i.e.,

$$271 \quad \mu_H = \frac{1}{2}\mu_{H_2} = \frac{1}{2}(\mu_{H_2,0} + RT \ln f_{H_2}) \quad (\text{A. 8})$$

272 where  $f_{H_2}$  is the fugacity. From Eqs. (A.1) and (A.8),

$$273 \quad \gamma_H = \frac{(f_{H_2})^{\frac{1}{2}}}{n_H} \exp \left\{ \frac{1}{RT} \left( \frac{1}{2}\mu_{H_2,0} - \mu_H^0 \right) \right\} \quad (\text{A. 9})$$

274 where it is noted that the stress gradient  $\nabla \sigma_h$  is 0 in the PCT measurement. Therefore,

$$275 \quad \frac{\partial \ln \gamma_H}{\partial \ln n_H} = \frac{n_H}{\gamma_H} \frac{\partial \gamma_H}{\partial n_H} = \frac{n_H}{2f_{H_2}} \frac{\partial f_{H_2}}{\partial n_H} - 1 \quad (\text{A. 10})$$

276 By substituting Eq. (A.10) into Eq. (A.7), Eq. (A.11) is obtained.

$$277 \quad D'' = D' \left( \frac{n_H}{2f_{H_2}} \frac{\partial f_{H_2}}{\partial n_H} \right) \quad (\text{A. 11})$$

278  $f_{H_2} = P_{H_2}$  at a few atm where the HSA actuator is used. The relationship between  $n_H$  and  
 279  $P_{H_2}$  is obtained from the PCT diagram [31].

280 Figure A.1 shows the variations of  $D'$  and  $D''$  with  $n_H$  at  $n_H < 0.03$  used in the present  
 281 work, assuming  $D = 3.75 \times 10^{-11}$  m<sup>2</sup>/s [35],  $V_H = 1.73$  cm<sup>3</sup>/mol [34], and  $V_M = 8.864$   
 282 cm<sup>3</sup>/mol [37].

## 283 **Appendix B. Simulation method**

284 To simulate the deformation of the actuator with hydrogen diffusion, a diffusion-  
 285 mechanical coupled analysis was performed using the finite element method. Governing  
 286 equations used were Eq. (1) and Eq. (B.1) in the diffusion and mechanical analysis,  
 287 respectively.

288

$$\boldsymbol{\sigma} = \boldsymbol{c}\boldsymbol{\varepsilon} + \boldsymbol{\beta}n_H \quad (\text{B.1})$$

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where  $\boldsymbol{c}$  is the stiffness tensor given by Young's modulus  $E$  and Poisson's ratio  $\nu$ ,  $\boldsymbol{\beta}$  is a tensor in which diagonal components are the linear expansion ratio,  $b$  induced by hydrogen insertion and non-diagonal components have zero values. The calculation procedure is shown in [28].

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Table B.1 Material properties used in analysis.

	$E$	$\nu$	$V_H$	$V_M$	$b$
	[GPa]	[-]	[cm <sup>3</sup> /mol]	[cm <sup>3</sup> /mol]	[-]
Palladium	121 [36]	0.39 [36]	1.73 [34]	8.864 [37]	0.0652
Epoxy resin [38]	3.80	0.25	-	-	-

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The distribution of  $n_H$  obtained from the diffusion analysis was input as a boundary condition in the mechanical analysis, and symmetrical boundary conditions were applied

309 on  $x = y = 0$ .

310 In the diffusion analysis, the stress distribution obtained from the mechanical analysis  
311 at the previous iteration was applied to each node.  $n_H$  at the surface layer part was  
312 determined as follows. Based on the reaction  $H_2(g) \Leftrightarrow H+H$  (solid solution), assume that  
313 rates of inflow  $J_{H_2,in}$  and outflow  $J_{H,out}$  on the foil surface are in proportion to the pressure  
314 and the square of the activity under free stress  $a_{H,0}$ , respectively. The surface flux of  
315 hydrogen atoms  $J_S$  is represented as

$$316 \quad J_S = 2J_{H_2,in} - J_{H,out} = 2k_{in}P_{H_2} - k_{out}a_{H,0}^2 \quad (B.2)$$

317 where  $k_{in}$  and  $k_{out}$  are constants. At the equilibrium condition,  $\mu_{H_2} = 2\mu_H$  and  $J_S = 0$ .  
318 Therefore,

$$319 \quad J_S = K_t(P_{H_2} - P_{H_2}^0 K_S^2 a_{H,0}^2) \quad (B.3)$$

320 where  $K_t = 2k_{in}$ . The activity  $a_H$  changes under stress, i.e.,

$$321 \quad a_H = a_{H,0} \exp\left(\frac{V_H \sigma_h}{RT}\right) \quad (B.4)$$

322 By substituting Eq. (B.4), Eq. (B.3) becomes

$$323 \quad J_S = K_t \left\{ P_{H_2} - P_{H_2}^0 K_S^2 a_H^2 \exp\left(-\frac{2V_H \sigma_h}{RT}\right) \right\} \quad (B.5)$$

324 The hydrogen content at the surface boundary layer at time  $t+1$  is represented as Eq. (B.6).

$$325 \quad n_H(t+1) = n_H(t) + \frac{K_t A_S}{V_S} \left\{ P_{H_2} - P_{H_2}^0 K_S^2 a_H^2 \exp\left(-\frac{2V_H \sigma_h}{RT}\right) \right\} \Delta t \quad (B.6)$$

326 where  $A_S$  is the surface area,  $V_S$  is the volume of the layer and  $\delta t$  is the time step. In this  
327 work,  $P_{H_2}$  and  $K_t$  were set as 1 atm and  $1 \times 10^{-4}$  mm/atm/s, respectively.

328 Figure B.2 shows a variation of the displacement  $U_{FEM}$  with time obtained from the  
329 numerical simulation, assuming that hydrogen of 1 atm was charged at  $t = 0$  s.  $U_{FEM}$   
330 decreased, i.e. the capsule shrank in the height direction, with the increase in time.  $u_{max}$   
331 was not shown because the effect of thermal expansion was ignored. The deformation

332 was completed after 3000 s, and the final displacement was -6.9  $\mu\text{m}$ . The response time  
333 was shorter than the experimental result. We assumed  $K_t = 1 \times 10^{-4}$  mm/atm/s, which is not  
334 a measured value and sensitive to the surface condition. This is why the deformation rate  
335 was different in the simulation and experiment. However, the difference is not discussed  
336 in this paper because the response time does not affect the internal state qualitatively.

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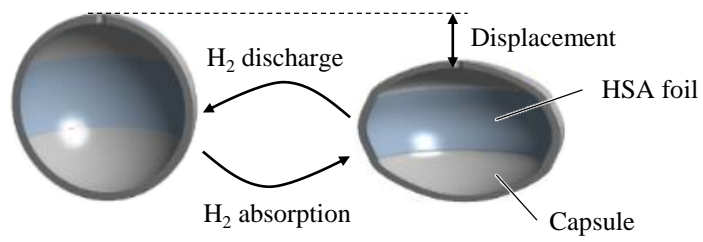
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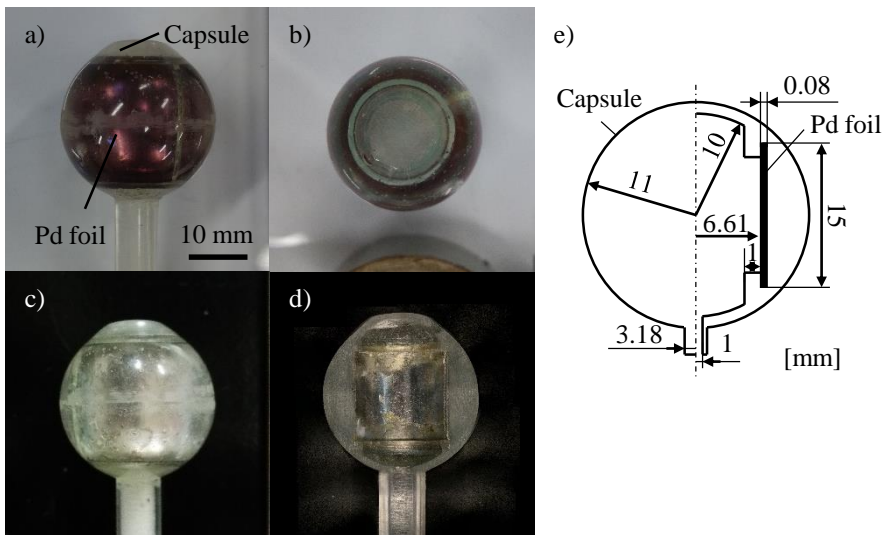
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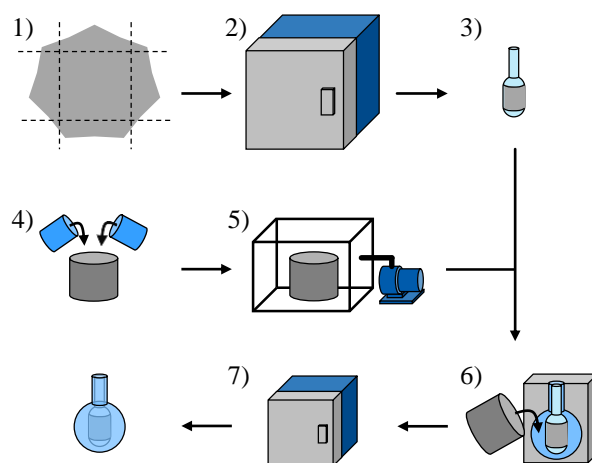
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**Figure 1** Schematics of capsule-type hydrogen-storage-alloy (HSA) actuator and its deformation.

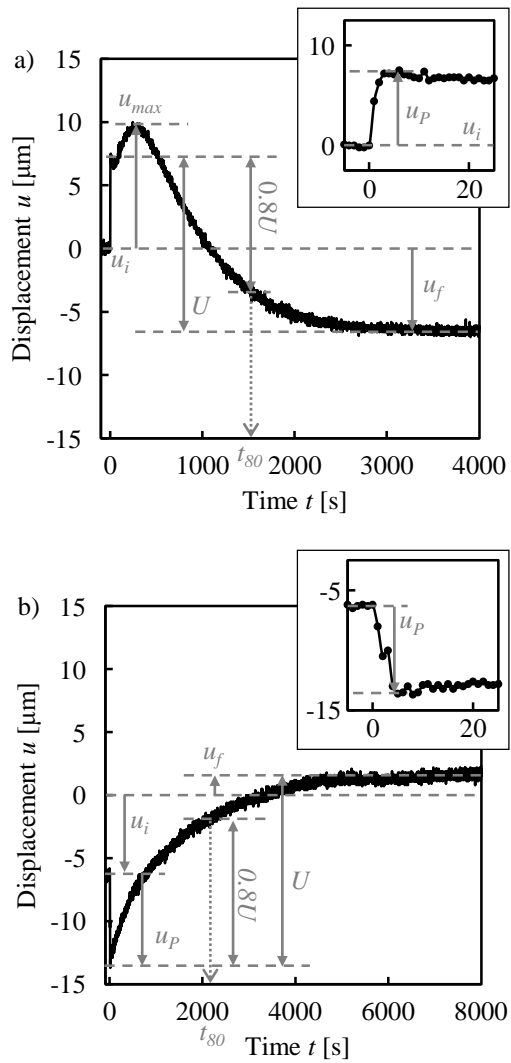


**Figure 2** Sample : a) side and b) top view of outer surface after fabrication, c) outer surface and d) cross section after experiment, and e) schematic image with design dimensions.

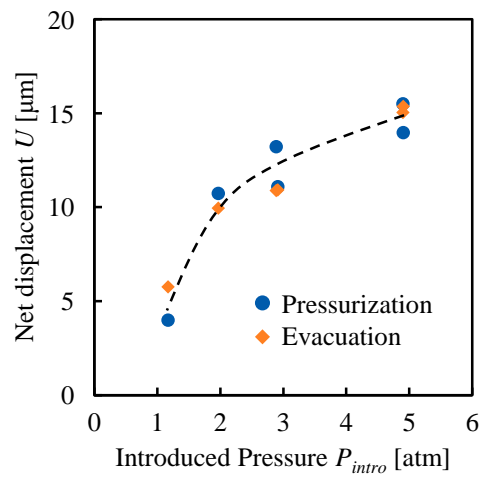


**Figure 3** Fabrication procedure of capsule-type HSA actuator: 1) cutting HSA foil, 2) annealing HSA foil, 3) winding HSA foil around core, 4) mixing Epon 828 and DETA, 5) vacuum degassing, 6) injection epoxy resin into mold with core, and 7) discharge of core by melting.

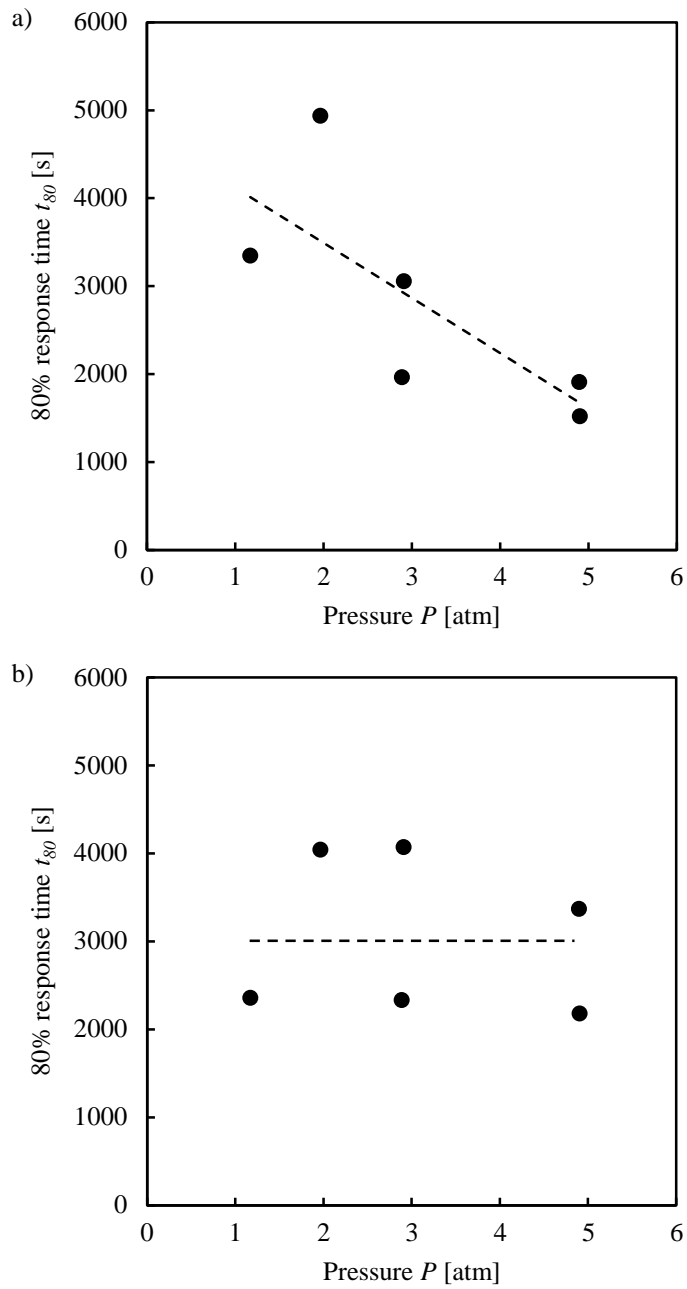




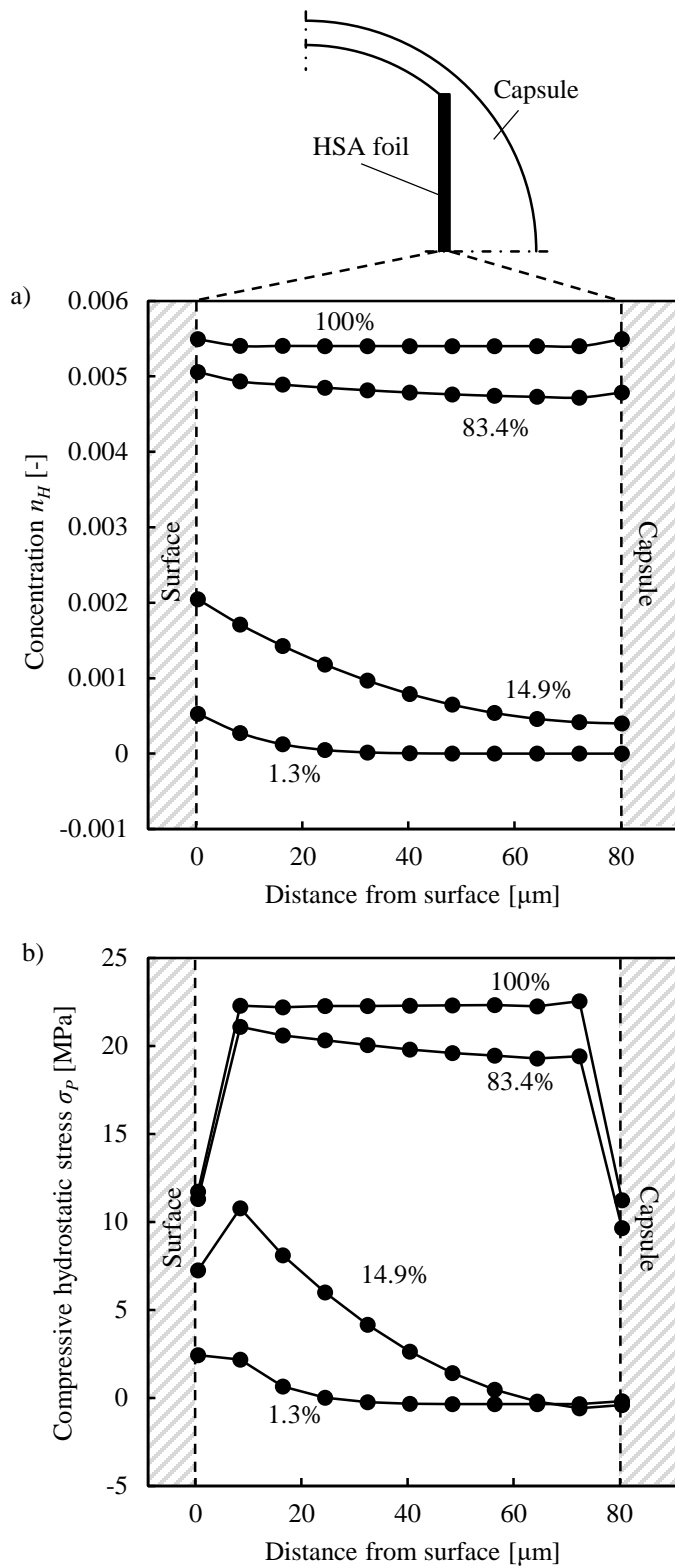
**Figure 4** Time change of displacement a) during pressurization and b) at following evacuation at  $P_{intro}$  of 4.903 atm.



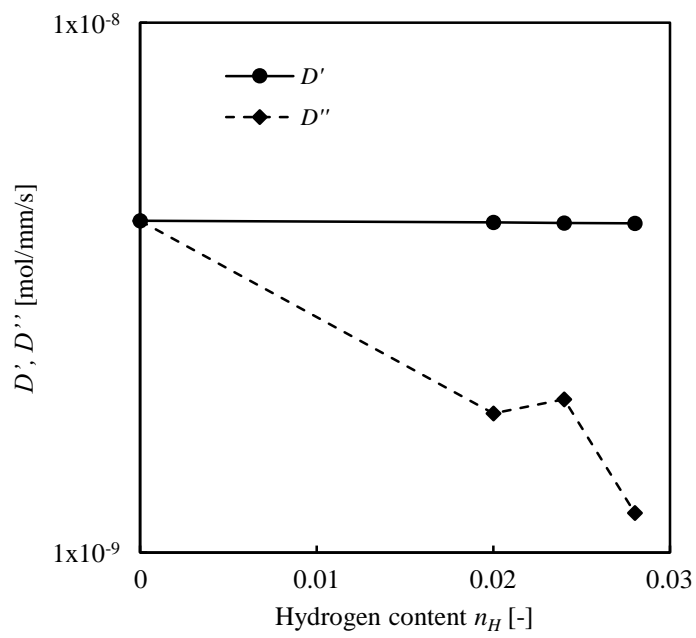
**Figure 5** Dependence of net displacement on introduced  $\text{H}_2$  pressure.



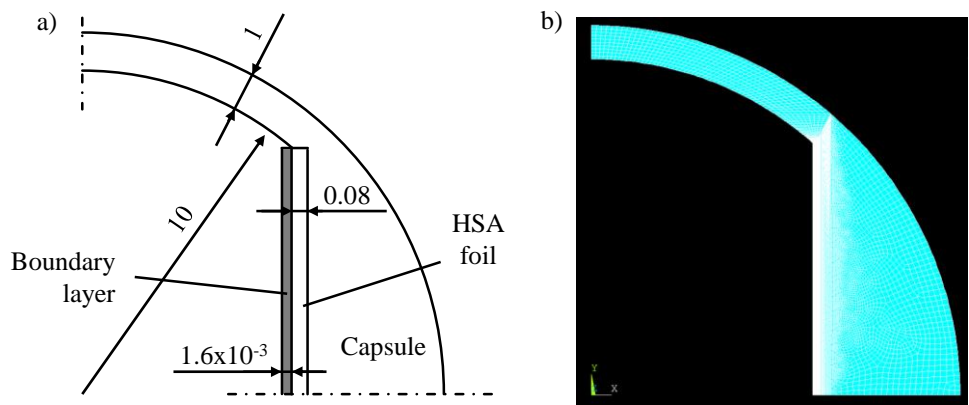
**Figure 6** Dependence of response time when deformation ratio reached 80% on pressure at a)  $H_2$  introduction and b) evacuation



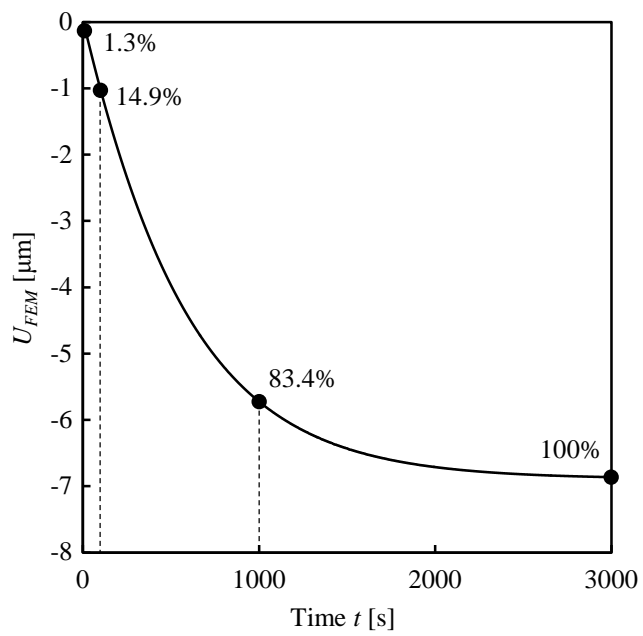
**Figure 7** Distribution of a) hydrogen content and b) compressive hydrostatic stress in hydrogen-storage-alloy foil at various deformation ratio. The deformation ratio was calculated by divided by the final displacement.



**Figure A.1** Variation of modified diffusion coefficients  $D'$  and  $D''$  with hydrogen content.



**Figure B.1** a) Schematic of analytical model and b) finite element model for diffusion-structural coupling analysis of capsule-type HSA actuator.



**Figure B.2** Displacement change with time at  $P_{intro} = 1$  atm in numerical simulation. It corresponds to doubled displacement at the top of the analytical model. The value at each plot represents the deformation ratio.